

Gravitational waves from proto-neutron star evolution

Giovanni Camelio

in collaboration with: Leonardo Gualtieri, Alessandro Lovato,
Jose A. Pons, Omar Benhar, Morgane Fortin & Valeria Ferrari

PhD student @ University of Rome "Sapienza"
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Supernovae (SNe)

- ▶ $8 M_{\odot} \lesssim M_{\star} \lesssim 25 M_{\odot} \rightarrow$ neutron star (NS);
- ▶ $E_{\gamma} = 10^{52}$ erg ($L_{\text{SN}} \simeq L_{\text{galaxy}}$);
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- ▶ Fe core supported by electrons Fermi pressure;
- ▶ when $M_{\text{core}} > M_{\text{Ch}} = 1.44 M_{\odot} \rightarrow$ core collapse;
- ▶ increasing density \rightarrow nucleons Fermi pressure $n_{\text{B}} \simeq n_0$;
- ▶ outer-core bounce \rightarrow shock wave;

Proto-neutron star (PNS)

PNSs are the SN contracting cores:

- ▶ very early evolution (PHASE I): core bounce $\div \sim 0.2$ s:
 - ▶ fully relativistic, highly dynamical codes;
 - ▶ mass accretion;
 - ▶ PNS contraction 150 km \rightarrow 30 km;
 - ▶ high-temperature PNS envelope;
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- ▶ early evolution (PHASE II): ~ 0.2 s \div minutes:
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 - ▶ **beta equilibrium**;
 - ▶ deleptonization stage (heating of the inner core);
 - ▶ cooling stage.

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- ▶ minutes: birth of a mature **neutron star** (neutrino transparent).

We are interested in the **gravitational wave emission** (from rotation or stellar oscillations) of the PNS in **PHASE II**.

PNS evolution: general facts

- ▶ in a “cold” NS the EoS is barotropic: $P \rightarrow (n_B, \epsilon, \dots)$
- ▶ in a PNS the $T \simeq 40 \text{ MeV} \gtrsim E_F \simeq 10 \text{ MeV}$ and therefore **the EoS is non-barotropic** $(s, Y_L, P) \rightarrow (n_B, \epsilon, T, Y_\nu, \dots)$

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- ▶ PNS structure from general relativistic TOV equations (spherical metric) with a given finite-temperature EoS $\epsilon(s, Y_L, P)$;
- ▶ ν transport (Boltzmann–Lindquist Eqs, **BLE**) with β -equilibrium to evolve the profiles of entropy s and lepton number Y_L ;

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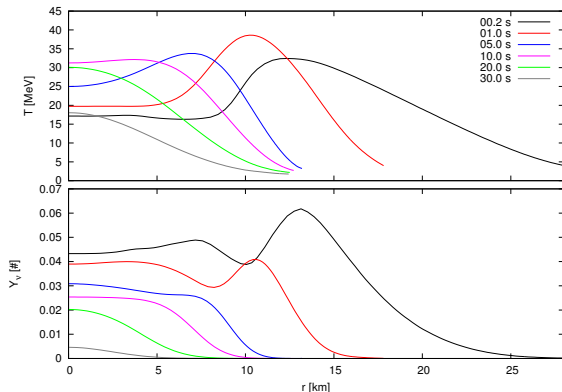
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- ▶ **neutrino diffusion coefficients** depend on the neutrino cross sections (and therefore on the EoS...);

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- ▶ **neutrino diffusion coefficients** depend on the neutrino cross sections (and therefore on the EoS...);
- ▶ For now, only mean-field EoSs have been used (e.g., GM3 Glendenning & Moszkowski, “Reconciliation of Neutron-Star Masses and Binding of the Λ in Hypernuclei”, PRL **67**:2414–2417 [1991]).

PNS evolution: our code

Our code reproduces the results of Pons, Reddy, Prakash, Lattimer & Miralles, "Evolution of proto-neutron stars", ApJ **513**:780–804 [1999]:



- ▶ low T core;
- ▶ high T envelope;
- ▶ trapped ν ;
- ▶ inner core heating;
- ▶ cooling;
- ▶ deleptonization.

Figure: PNS evolution, GM3 EoS (our code).

Effective inclusion of rotation: the procedure

We have extended the work of Villain, Pons, Cerdá-Durán & Gourgoulhon, “Evolutionary sequences of rotating protoneutron stars”, *A&A* **418**:283–294 [2004]:
Camelio, Gualtieri, Pons & Ferrari, “Spin evolution of a proto-neutron star”, *PRD* **94**, 024008 (2016), [arXiv:1601.02945](https://arxiv.org/abs/1601.02945) [astro-ph.HE].

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First, evolve the non-rotating star:

- ▶ fix the total baryon mass M_b ;
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- ▶ the (non-rotating) evolution gives $s(t, a)$, $Y_L(t, a)$.

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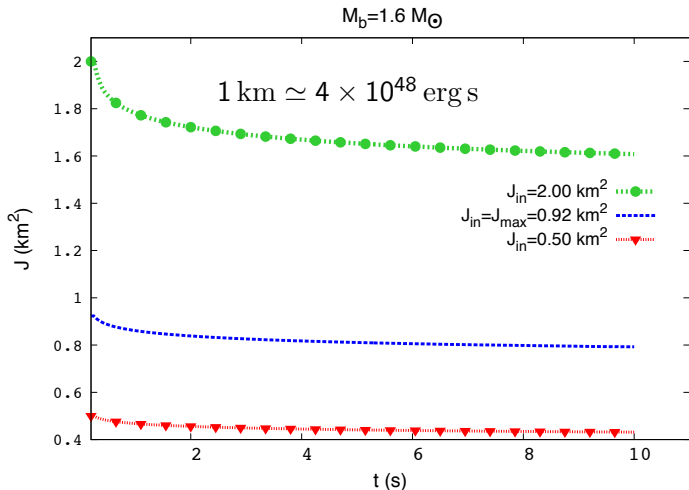
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To effectively include the rotation:

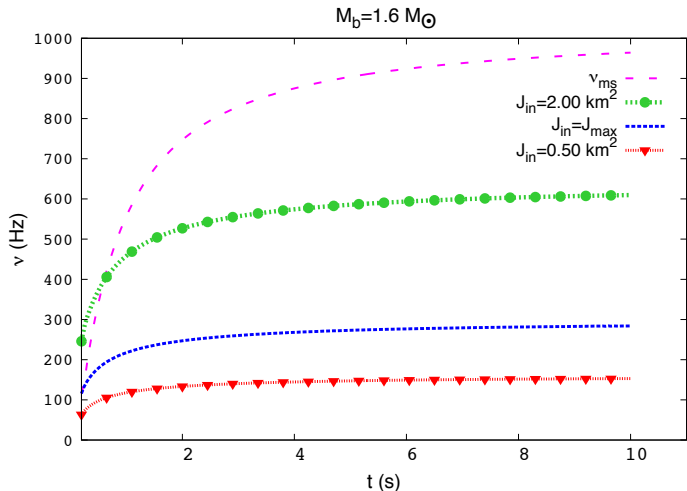
1. fix an initial angular momentum $J_{in} = J(t = 0)$;
2. “effective” EoS at time t : $\epsilon'_t(a, P) = \epsilon(s(t, a), Y_L(t, a), P)$;
3. using ϵ'_t , solve Hartle-Torne (structure equations of a slowly rigidly rotating PNS) at time t with fixed M_b and $J(t)$;
4. determine $J(t + dt)$ using the Epstein formula;
5. $t \rightarrow t + dt$, back to point 2.

Effective inclusion of rotation: results



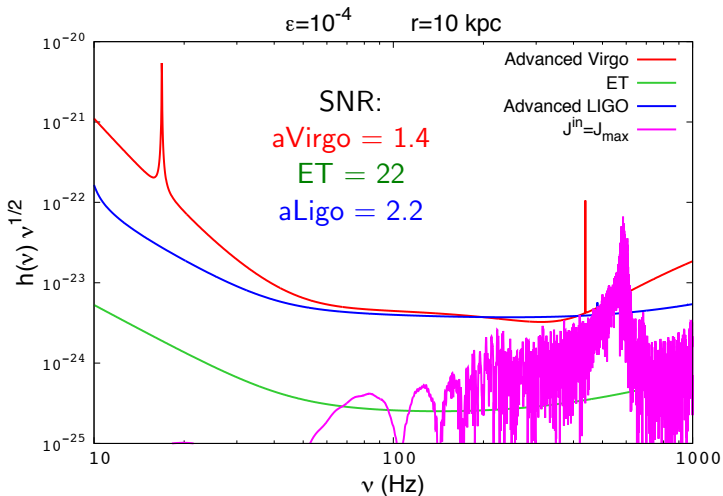
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EoS dependence: general facts

- ▶ previous PNS evolution studies used relativistic mean-field EoSs; we want to use a general EoS, e.g. a nuclear many-body theory EoS (Lovato & Benhar, “An effective interaction from Argonne-Urbana nuclear forces”, in preparation);

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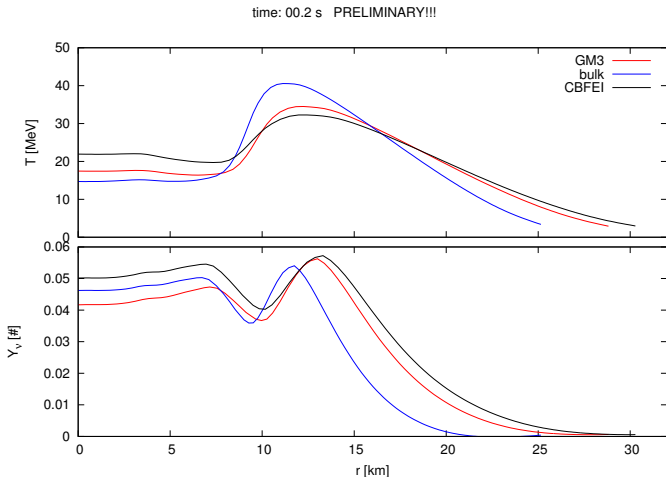
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- ▶ from the free energy per baryon $f(T, n_b, Y_p)$ you can obtain all the other thermodynamical quantities with derivatives!
- ▶ we want to obtain the EoS from the fit of the interacting part of the baryon free energy

$$\begin{aligned}f_{\text{EoS}}(T, n_b, Y_p) &= f_{\text{free gas}}(T, n_b, Y_p) + f_{\text{I}}(T, n_b, Y_p), \\f_{\text{I}}(T, n_b, Y_p) &= 4Y_p(1 - Y_p)f_{\text{SNM}}(T, n_b) \\&\quad + (1 - 2Y_p)^2 f_{\text{PNM}}(T, n_b), \\f_{*\text{NM}}(T, n_b) &= \text{polynomial in } T \text{ and } n_b,\end{aligned}$$

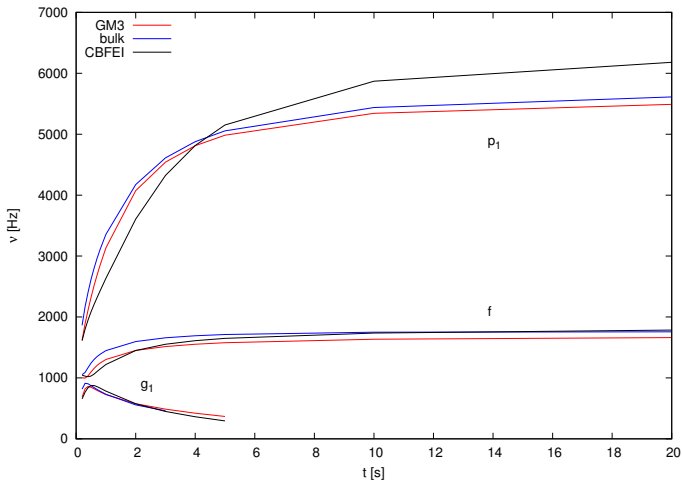
that is similar to how the bulk nuclear matter has been treated in Lattimer & Swesty, “A generalized equation of state for hot, dense matter”, Nucl.Phys.A **535**:331 [1991].

EoS dependence: results (preliminary!)



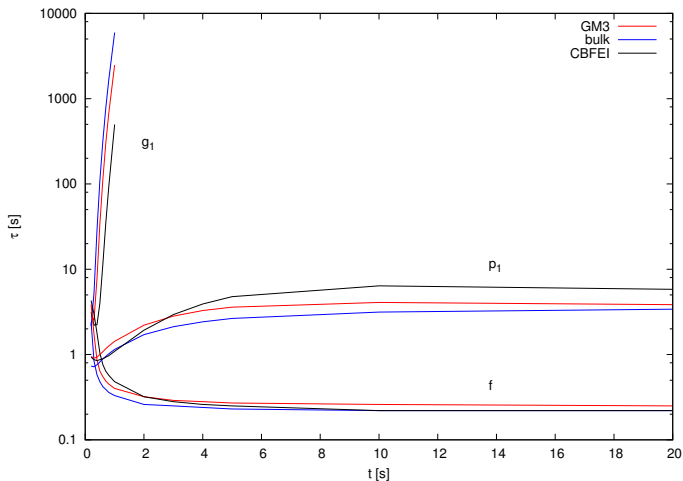
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Done:

- ▶ new PNS evolution code;
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Outlooks:

- ▶ convection (mixing length theory);
- ▶ accretion;
- ▶ evolution in 1+1.5D (consistent inclusion of rotation).