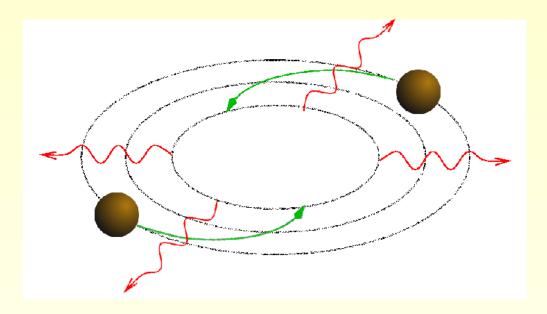
Evolution with the conformal field equations

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Introduction



- study a self–gravitating system
 - formation of singularities, horizons
 - extract radiation
 - global structure
- modelled as isolated system
- space–time is asymptotically flat \iff can be conformally compactified

Conformal compactification of Minkowski space-time

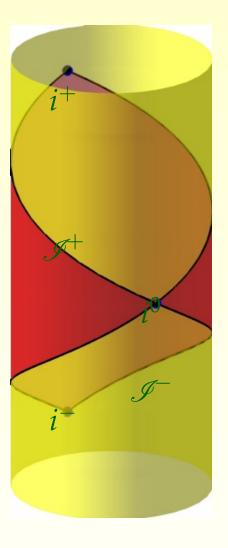
$$\tilde{g} = dt^2 - dr^2 - r^2 d\omega^2, \qquad d\omega^2 = \left(d\theta^2 + \sin^2\theta d\phi^2\right)$$
$$= \frac{1}{(\cos T + \cos R)^2} \left(dT^2 - dR^2 - \sin^2 R \, d\omega^2\right),$$

has the form

$$\tilde{g} = \Omega^{-2}g \iff g = \Omega^2 \tilde{g}.$$

- \mathscr{I}^+ and \mathscr{I}^- form a 3–D boundary
- points on \mathscr{I} are at infinity wrt \tilde{g}
- conformal factor $\Omega = \cos T + \cos R = 0$ on \mathscr{I}
- physical space-time (*M̃*, *g̃*) is smoothly embedded into a larger space-time (*M*, *g*) with the same light-cone structure.

 \rightarrow \mathscr{I} is an abstract concept which allows us to talk about radiation



Conformal field equations

- express $\tilde{G}_{ab} = 0$ in terms of $(\mathcal{M}, \Omega, g_{ab})$
- conformal transformation

$$0 = \tilde{G}_{ab} = G_{ab} - \frac{2}{\Omega} \left(\nabla_a \nabla_b \Omega - g_{ab} \Box \Omega \right) - \frac{3}{\Omega^2} g_{ab} \nabla_c \Omega \nabla^c \Omega.$$

• read as equation for Ω ($G_{ab} = -2\Phi_{ab} - 6\Lambda g_{ab}$)

$$\nabla_a \nabla_b \Omega = \Omega \Phi_{ab} + g_{ab} S, \qquad S = \frac{1}{4} \Box \Omega$$

• need equation for $\Phi_{ab} \rightarrow \text{Bianchi}$ identities in M and \tilde{M}

$$\tilde{\nabla}_{d}C_{abc}{}^{d} = 2\nabla_{[a}\tilde{\Phi}_{b]c} + 2\tilde{g}_{c[a}\nabla_{b]}\tilde{\Lambda} = 0 \iff \nabla_{a}\underbrace{\left(\Omega^{-1}C_{abc}{}^{d}\right)}_{K_{abc}{}^{d}} = 0$$

 $\nabla_d C_{abc}{}^d = 2\nabla_{[a} \Phi_{b]c} + 2g_{c[a} \nabla_{b]} \Lambda \iff 2\nabla_{[a} \Phi_{b]c} = \nabla_d \Omega K_{abc}{}^d - 2g_{c[a} \nabla_{b]} \Lambda$

Conformal field equations

$$\Gamma = \partial g \qquad \qquad \partial e = \gamma e \partial^2 g + \partial g \partial g = \Omega K + \Phi + \Lambda \qquad \qquad \partial \gamma + \gamma \gamma = \Omega K + \Phi + \Lambda \partial \Phi = \partial \Omega K - g \partial \Lambda \partial K = 0 \partial^2 \Omega = \Omega \Phi + g S \partial S = \Phi \partial \Omega + \Omega \partial \Lambda + \Lambda \partial \Omega$$

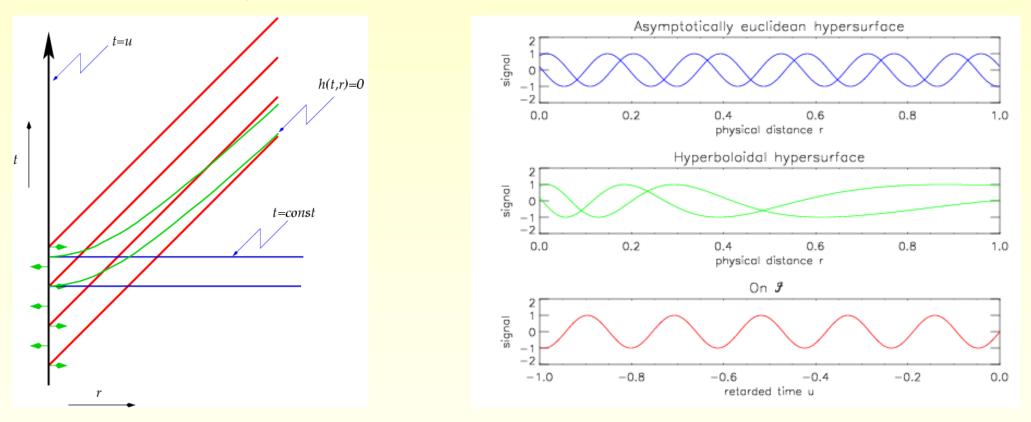
gauge source functions

$$\Box x^{\mu} = F^{\mu} \qquad \Box x^{\mu} = F^{\mu}$$
$$\nabla_{a}(e_{i}\nabla^{a}e_{k}) = F_{ik}$$
$$\Lambda \qquad \Lambda$$

Analysis

Friedrich (1979-)

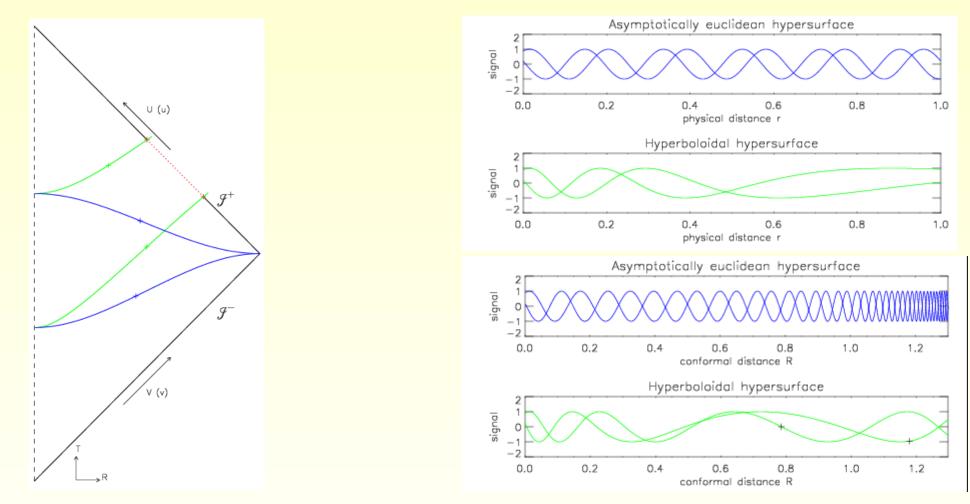
- Split into constraints and symmetric hyperbolic evolution equations
- evolution propagates the constraints
- hyperboloidal IVP is well posed
- (semi)–global existence for small hyperboloidal data
- gluing techniques give global existence for asymptotically flat manifolds
- initial data on hyperboloidal hypersurfaces can be constructed



Waveforms on hypersurfaces in Minkowski space

- $t = t_0$ -surface 'collects' all waves emitted up to $t = t_0$
- hyp. surface 'selects' waves emitted in a finite interval before $t = t_0$
- $t = t_0$ -surface is 'instantaneous', hyp. surface 'slows down'
- spatial succession turns into temporal succession

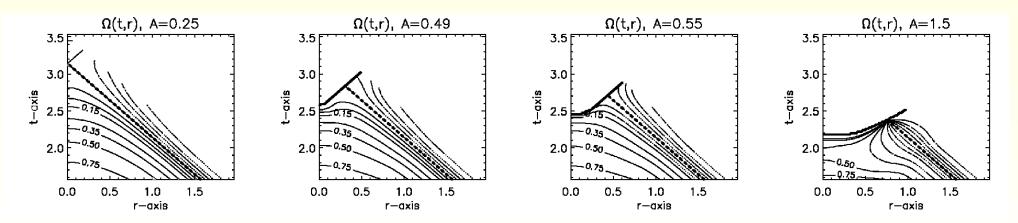
Waveforms on hypersurfaces in conformal Minkowski space



- piling up of oscillations on asymptotically euclidean hsf 'all waves at once'
- smooth across null–infinity

P. Hübner (1994)

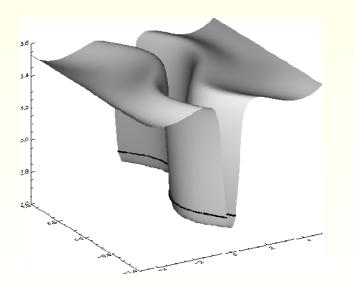
- spherical symmetry with conformally coupled scalar field
- parameter study of collapse
- numerical generation of conformal diagrams
- investigation of the causal nature of the singularity

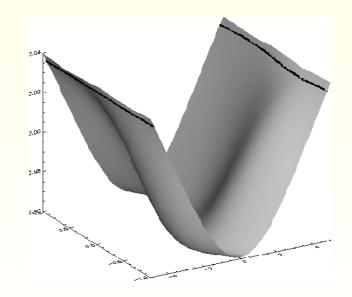


JF (1998)

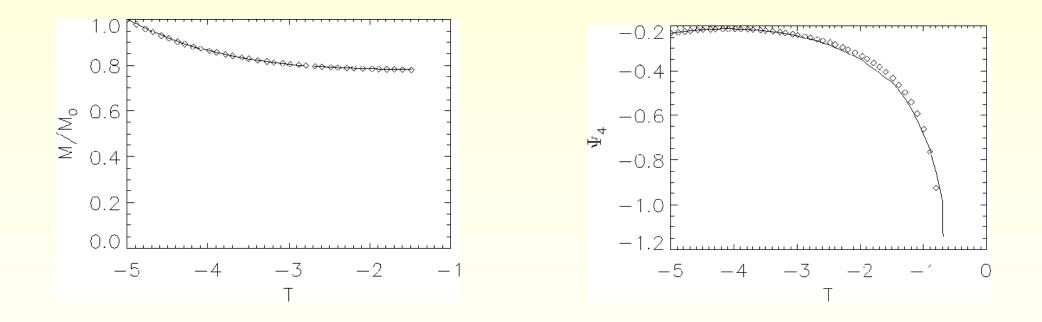
- frame based vacuum code in 2 + 1 dimensions
- space–like non–vanishing Killing vector with closed orbits: 'toroidal' \mathscr{I}
- 'generalized harmonic' time coordinate $\Box t = F$, shift components free
- Fermi transport of the frame
- Investigate radiation extraction

Several gauge choices, 'scri-freezing'





radiation extraction on \mathscr{I}

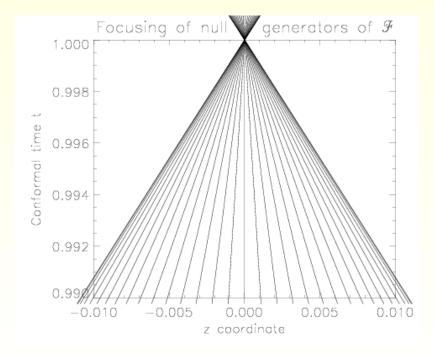


JF (1998)

- spectral solver for hyperboloidal initial data following Andersson, Chrusciel, Friedrich
- choose conformal 3-metric *h*_{ab}
- solve Yamabe equation for conformal factor so that $\tilde{h}_{ab} = \Omega^{-2} h_{ab}$ has ${}^3\tilde{R} = -6$
- all other initial data can be computed
- draw–backs: numerical differentiation, 0/0-terms, nasty

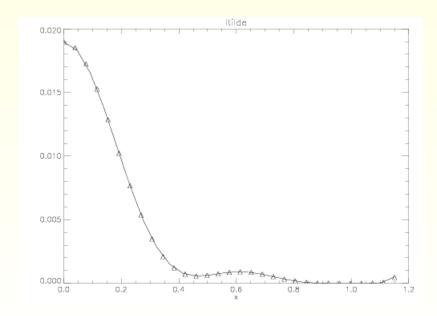
P. Hübner (2000)

- Metric based 3+1 code, no symmetries
- MOL with 4th order space discretisation and 4th order Runge–Kutta
- Outer boundary outside *I* treated by transition zone to static state
- gauge sources: densitised lapse, shift and scalar curvature Λ
- numerically integrated generators of \mathscr{I}



S. Husa (2002)

- Further development with Hübner's code
- Parameter studies
- Variations in the gauge sources
- weak Brill data compared with perturbative analysis
- problems with transition zone



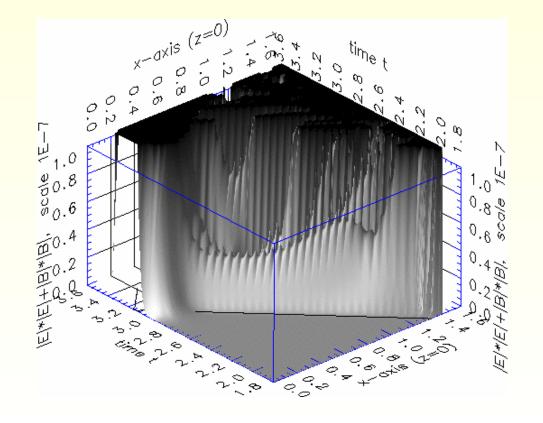
JF and M. Hein (2003)

- Frame based axi-symmetric code
- MOL with RK4 and space–discretisation using Cartoon method
- Outer boundary treated by eigen-field method
- gauge sources: F^{μ} , F_{ik} , Λ

Effect of *I* as a one-way membrane

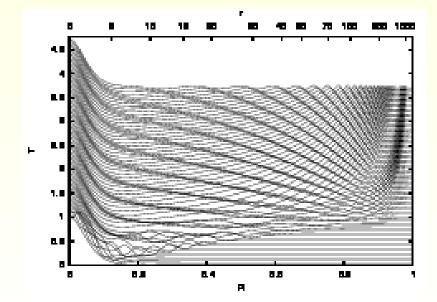
Minkowski space-time with white noise as a boundary condition

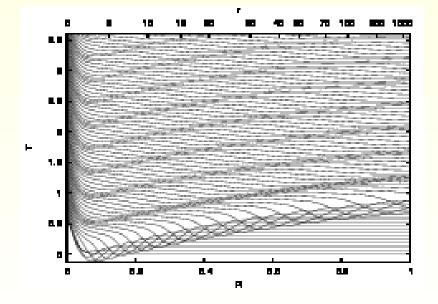
norm of the Weyl tensor



Fodor and Rácz (2005)

- Yang–Mills–Higgs system in spherical symmetry
- spherically symmetric perturbation of BPS monopole
- Conformal compactification and hperboloidal slicings
- breathing solution found
- good interaction with analytical community (Forgács, Volkov)





Other developments

- Development of 3+1 code within the Cactus framework
- Extension of Friedrich–Nagy boundary treatment to CFE
- Thin–sandwich approach for hyperboloidal initial data
- Initial data close to Kruskal data
- Perturbative analysis of constraint propagation for the Weyl system
- Preliminary work on the General CFE
- Exploration of conformal Gauss gauge on Schwarzschild and Kerr
- Need a way to solve the conformal constraints directly!