

7 Heisenberg's Uncertainty Principle

In this section we again consider a particle moving in one dimension only. We have seen that a wavepacket of matter waves describes a particle which is not located at a particular point. Note that this is not the same as saying that its position is unknown, which implies - falsely - that its exact position exists but that the experimentalist does not know what it is. The particle simply does not have a well-defined position until a measurement is made at which point the wavefunction of the particle is changed into one representing a particle of well-defined position (a very narrow wavepacket). Instead we only know what the probability is that the result of such a measurement of its position will yield a given result. This probability is very small, except in a region of width x , and so we identify this range, Δx , as the “uncertainty” in the position of the particle.

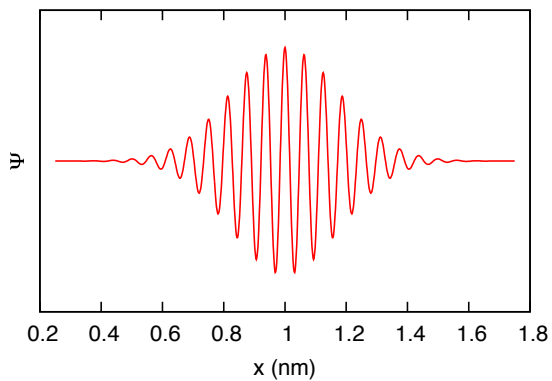
We have seen when discussing wavepackets, that a wavepacket of a given width is constructed from a superposition of waves with wavelength, λ , and amplitude $A(\lambda)$. The amplitude $A(\lambda)$ has an interpretation which is analogous to the interpretation of the amplitude, $|\Psi(x, t)|$, of the wave in the wavepacket, namely that its square, $|A(\lambda)|^2 d\lambda$, is the probability that a measurement of the wavelength of the particle will yield a value between λ and $\lambda + d\lambda$. By the de Broglie wave relation, the wavelength measurement is equivalent to a measurement of the momentum. Since the momentum is actually inversely proportional to λ (i.e. proportional to $1/\lambda$),⁹ it is more convenient to plot $A(\lambda)$ as a function of $1/\lambda$. We have a distribution whose square tells us how likely it is that a measurement of $1/\lambda$, will yield a given result and we can see that such a probability is small except over a small range of $1/\lambda$. The particle does not possess a single wavelength, but is a superposition of waves of different wavelengths with different amplitudes. The range over which this amplitude is not small represents the uncertainty in $1/\lambda$.

We recall that the width of the wavepacket, Δx , is inverse to the width of the distribution of $A(\lambda)$ whose width is $\Delta(1/\lambda)$, i.e for a wavepacket with large width, Δx , representing a particle with a large uncertainty in its position, the distribution $A(\lambda)$ is narrow implying a small uncertainty in the quantity $1/\lambda$. On the other hand for a wavepacket with small width, Δx , representing a particle with a small uncertainty in its position, the distribution $A(\lambda)$ is wide implying a large uncertainty in the quantity $1/\lambda$.

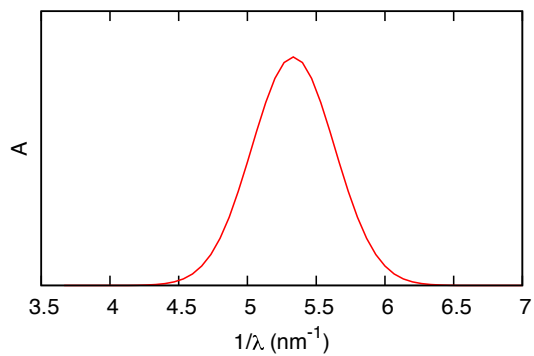
This is demonstrated in Fig. 32. The graphs on the left are wavepackets. The upper wavepacket has a width of about 0.8 nm whereas the lower graph is a narrower wavepacket with a width of 0.4 nm. The graphs on the right are the distributions of the amplitude at wavelength, λ , plotted against $1/\lambda$. The upper distribution has a width of about 1.2 per nm whereas the lower distribution has a width of about 2.4 per nm. In both cases the product of the width in position, x , and the width in the inverse wavelength, $\Delta(1/\lambda)$ is approximately one.

From the de Broglie wave relation, the uncertainty, Δp , in the momentum is h (Planck's

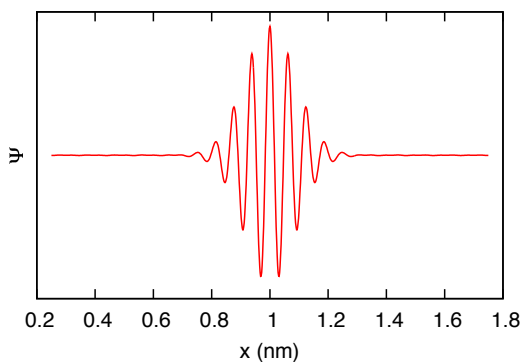
⁹The inverse of the wavelength of a wave is called its **wavenumber** and is often denoted by k but we will maintain the notation $1/\lambda$.



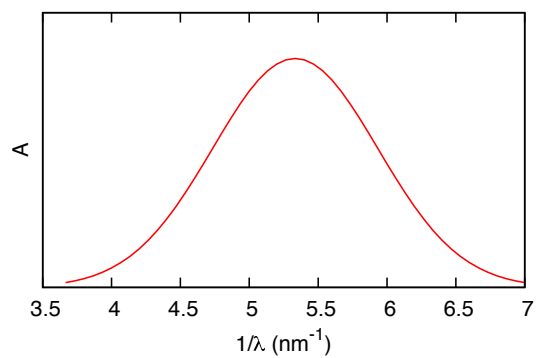
(a)



(b)



(c)



(d)

Figure 32: (a): a wavepacket with a width of 0.8 nm; (b): the corresponding distribution in amplitude of wavelength λ (plotted against $1/\lambda$), with a width of 1.2 per nm. (c): a wavepacket with a narrower width of 0.4 nm; (d): the corresponding distribution in amplitude of wavelength λ (plotted against $1/\lambda$), with a broader width of 2.4 per nm.

constant) times $\Delta(1/\lambda)$, so that if the product of Δx and $\Delta(1/\lambda)$ is approximately one, then the product of Δx and Δp is approximately h and we have the Heisenberg uncertainty relation [24]

$$\Delta x \times \Delta p \approx h \tag{7.1}$$

The particle, whose matter wave is a wavepacket, does not have a well-defined (i.e. exact localized) position until its position is measured to within a given accuracy - at which point the wavefunction changes into a narrow wavepacket. Similarly, it is also the case that the particle whose wavefunction is described by that particular wavepacket, does not possess a well-defined momentum (since it is constructed out of a range of wavelengths) until its momentum is measured to a given accuracy,¹⁰ at which point the wavefunction changes to a much broader wavepacket but with a much narrower distribution in wavelength and hence a narrower distribution in momentum. One can measure the position accurately or the momentum accurately but not both. The product of the uncertainty, Δx , in the measurement of position (the width of the wavepacket) and the uncertainty, and Δp , in the measurement of momentum must be larger than Planck's constant, h , in keeping with the uncertainty relation.

To give you some idea of why we do not experience this uncertainty in our everyday observations, suppose we were looking under a microscope at a human blood cell whose mass is of order of a millionth of a microgram and the uncertainty in position is around one micron (a millionth of a meter - about the smallest resolution possible with visible light) then the uncertainty in its velocity due to Heisenberg's uncertainty relation would be of the order of one micron per year !!

¹⁰The momentum of a particle is proportional to its velocity (neglecting relativistic corrections) so that a measurement of a particles momentum is equivalent to a measurement of its velocity.