

Current-driven domain wall switching in multilayer exchange spring nanopillars

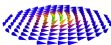
M. Franchin¹, T. Fischbacher¹, A. Knittel¹, P.A.J. de Groot²,
and H. Fangohr¹

¹ School of Engineering Sciences, University of Southampton (UK)

² School of Physics and Astronomy, University of Southampton (UK)

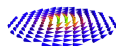
e-mail: franchin@soton.ac.uk

CMMP09, 16-12-2009



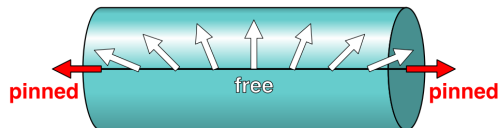
Idea: a domain wall responds “better”
to the application of a current
if it can rotate around its axis

- a novel domain wall (DW) structure: **rotating DW**
- **the physics** of a rotating DW
- **analytical** considerations
- **pushing** the DW through a potential barrier

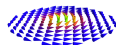


The system: rotating DW

In this talk we call **rotating domain wall** a Néel domain wall which is *constrained* in space, but *free to rotate* around its axis.



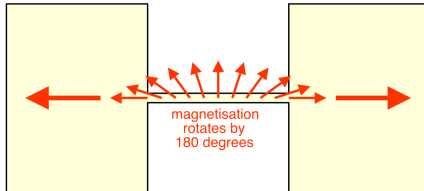
- ferromagnetic cylindrical nanopillar;
- magnetisation pinned along opposite directions at opposite faces of the cylinder;
- energy degeneracy: magnetisation free to rotate;
- inducing rotation does not “cost” any energy!



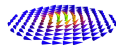
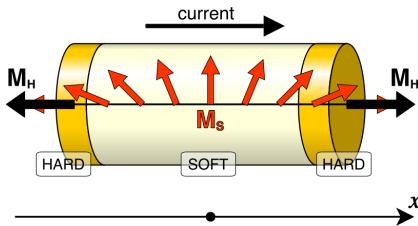
Possible realisations

Possible ways of constraining a domain wall:

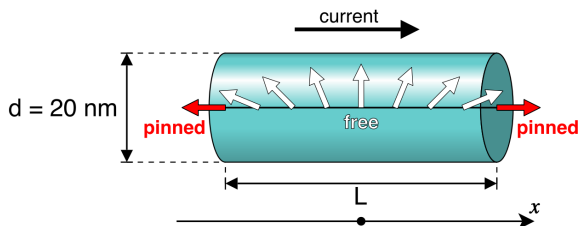
(a) geometrically



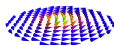
(b) exchange spring



Micromagnetic simulations



- Finite Element discretisation (Nmag, <http://nmag.soton.ac.uk>)
- Rigid pinning at opposite faces of cylinder ($\frac{d\vec{M}}{dt} = 0$)
- Applying current $j_P = Pj = 10^{11} \text{ A/m}^2$ along the axis (positive x direction)
- zero applied field, $\vec{H}_{\text{app}} = 0$
- Neglecting Oersted field and heating



Spin-currents in simulations

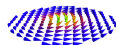
Zhang-Li correction to the LL equation (PRL 98, 187202 (2007)):

$$\begin{aligned}\frac{d\vec{M}}{dt} = & -\gamma \vec{M} \times \vec{H} + \frac{\alpha}{M_s} \vec{M} \times \frac{d\vec{M}}{dt} \\ & - \frac{v}{M_s^2} \vec{M} \times \left(\vec{M} \times \frac{d\vec{M}}{dx} \right) - \frac{\xi v}{M_s} \vec{M} \times \frac{d\vec{M}}{dx}\end{aligned}$$

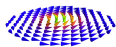
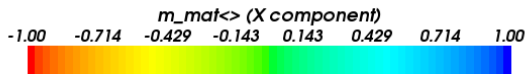
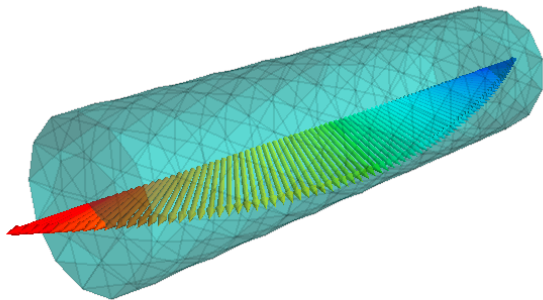
where,

$$v = \frac{Pj\mu_B}{eM_s(1 + \xi^2)},$$

P is the spin-polarisation and ξ a dimensionless number, expressing the “degree” of non-adiabaticity.



Animation



Precession frequency as a function of time

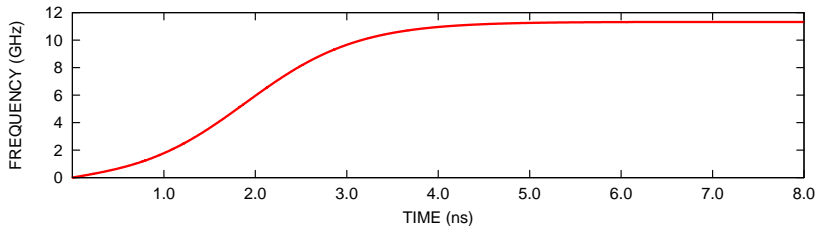
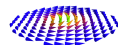


Figure: Time-evolution of the precession frequency.

The frequency of rotation of $\langle \vec{M} \rangle$ around the nanopillar axis increases monotonically toward an **asymptotic value**.

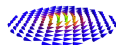


Asymptotic frequency as function of j and L

Result:

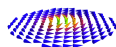
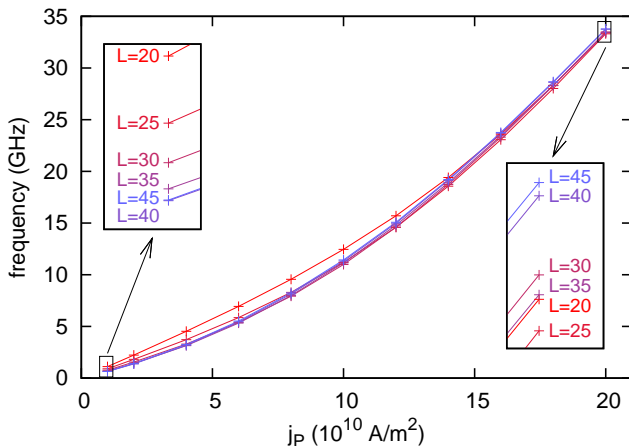
- Direct current (DC) gives rise to stationary precession!
- Nano-oscillator which requires no applied field!
- Could be used as a microwave generator!

How does the asymptotic frequency depend on the current density, j , and the nanopillar length, L ?



Asymptotic frequency as function of j and L

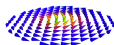
We carry out several simulations for different values of j and L :



One dimensional analytical model [PRB 78, 054447 (2008)]:
frequency ν as a function of current density j and damping α .

A critical current j_c can be identified, $j_c = \frac{2e\gamma}{\mu_0\mu_B} \frac{\alpha A}{L}$, and,

- LINEAR REGIME: $\nu \propto \frac{j}{\alpha L}$ for $j \ll j_c$
 - **no deformation** of DW
 - precession frequency increases **linearly** with the current
- QUADRATIC REGIME: $\nu \propto \left(\frac{j}{\alpha}\right)^2$ for $j \gg j_c$
 - **DW deformed** by current (gets compressed)
 - precession frequency increases **quadratically** with current



Validation of the analytical model

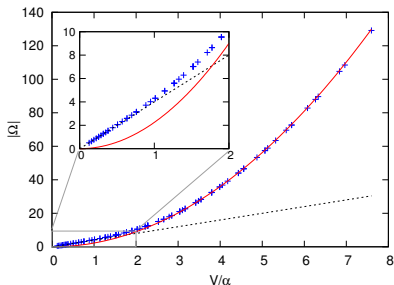


Figure: Analytical model (solid and dashed lines) and 1D micromagnetic simulations (blue crosses).

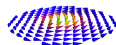
We obtained formulas for frequency as function of applied current j and nanopillar length L :

$$\nu = \begin{cases} \frac{2}{\pi^2} \frac{v}{\alpha L} & \text{for } j \ll j_c \\ \frac{1}{2\pi\gamma C} \left(\frac{3}{2} \frac{v}{\alpha}\right)^2 & \text{for } j \gg j_c \end{cases}$$

Recall that:

$$v = \frac{\mu_B P j}{e M_s (1 + \xi^2)} \propto j.$$

NOTE: j always appears as j/α

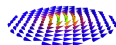


Energy pumping in the quadratic regime

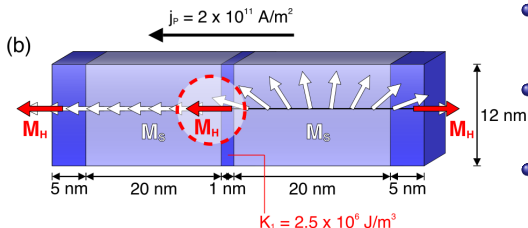
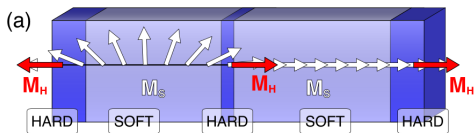
In the **quadratic regime**, the DW is compressed
⇒ **energy is pumped** into the DW!

IDEA:

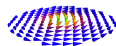
Use such energy to overcome an energy barrier:
push the domain wall through a hard layer!



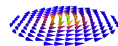
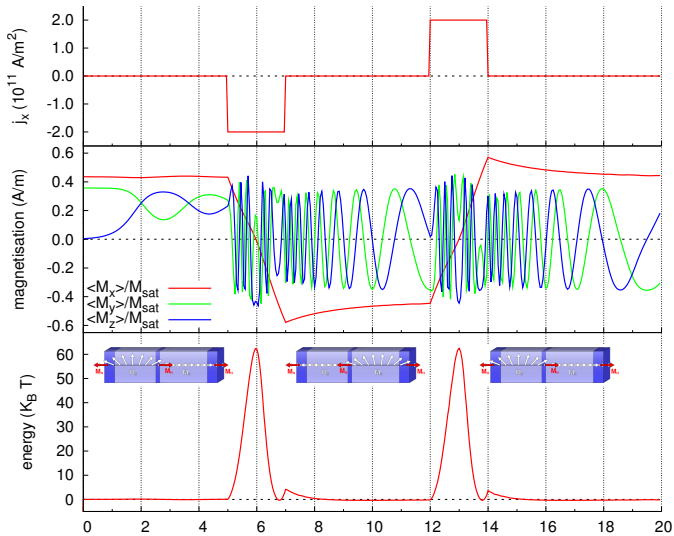
Current-driven switching of DW



- five-layer nanopillar
- square cross section (12 nm)
- “rigid” pinning in external layers
- “weak” pinning in central layer
- two stable states: DW on the left, DW on the right
- current to switch between the two states



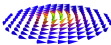
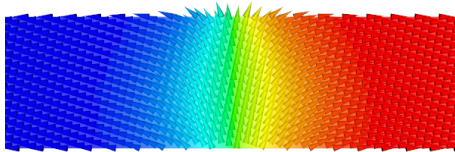
Current-driven switching of DW



(energy K_B T for $T = 300$ K)

Doubt:

Is this the usual current-driven domain wall motion?

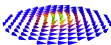
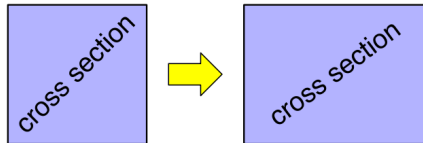


Doubt:

Is the rotation important at all?

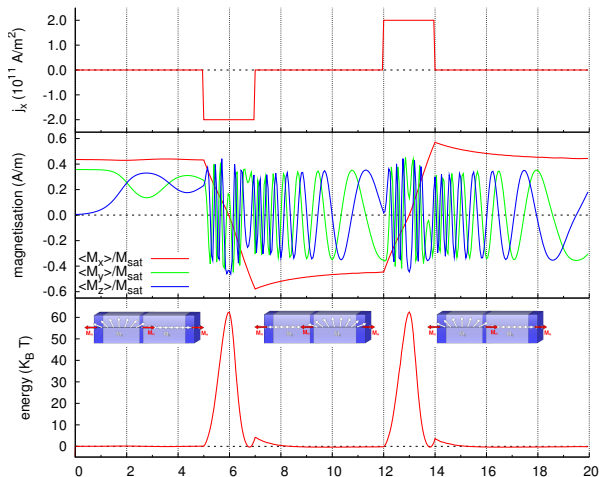
What happens if the DW is not free to rotate?

For example, let's choose a rectangular cross section...

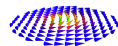


What happens with a rectangular cross section?

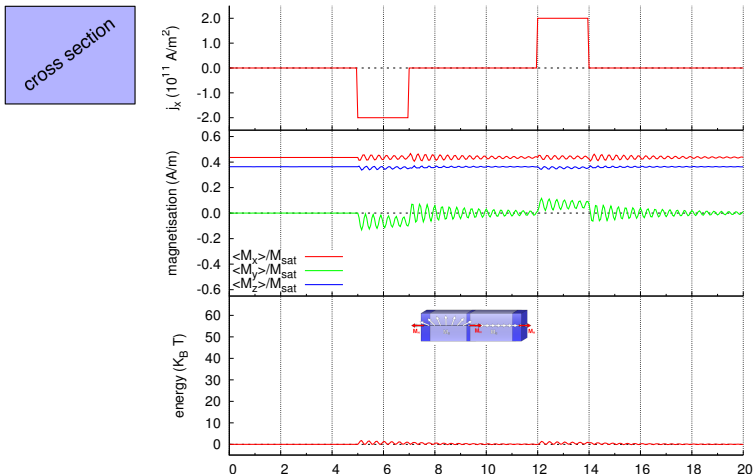
Cross section



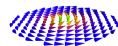
The DW is **not free to rotate**: nothing happens then!



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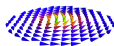


Conclusion:

A rotating DW can do things that a conventional DW can not!

It can:

- get in a **stationary precession** motion (microwave generation)
- **absorb** relevant quantities of **energy** (compression)

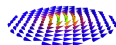


Summary

- A rotating domain wall can be used for **microwave generation** and
- Can be **pushed** effectively through a potential barrier
- New **switching** mechanism for MRAM?
- For more info, see:
 - M. Franchin *et al.*, J. Appl. Phys. 103, 07A504 (2008)
 - M. Franchin *et al.*, Phys. Rev. B 78, 054447 (2008)

Acknowledgement: The research leading to these results has received funding from the European Community's Seventh Framework Programme (FP7/2007-2013) under Grant Agreement n° 233552, and from EPSRC (EP/E040063/1)

Thank you!



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