

CHEM3023: Spins, Atoms and Molecules

Lecture 3

The Born-Oppenheimer approximation

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Learning outcomes

- Separate molecular Hamiltonians to electronic and nuclear parts according to the Born-Oppenheimer approximation
- Be able to manipulate expressions involving electronic wavefunctions, taking into account spin and space coordinates and antisymmetry

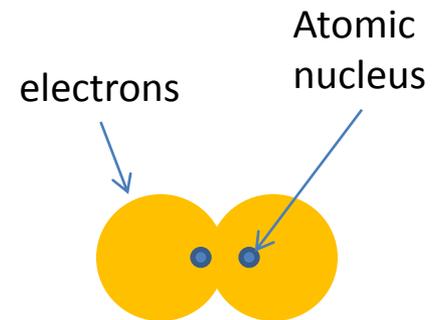
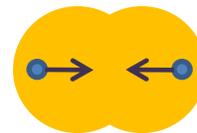
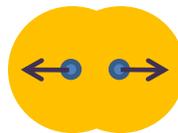
Separating electronic from nuclear coordinates

Nuclei are much heavier than electrons (e.g. Proton mass ≈ 1800 times the mass of an electron \Rightarrow compare kinetic energy operators of proton and electron).

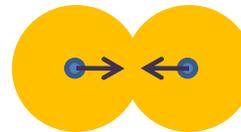
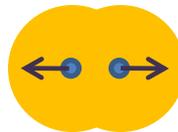
As a result nuclei move much slower than electrons. To a very good approximation, we can assume that the electronic motions are instantaneously “equilibrated” at each nuclear position.

Example: Vibration of diatomic molecule:

Wrong picture: “electron cloud” trails behind moving nuclei



Correct picture: “electron cloud” instantaneously re-arranges itself around moving nuclei



Hamiltonian operator for water molecule

Water contains 10 electrons and 3 nuclei. We will use the symbols "O" for the oxygen (atomic number $Z_O=8$) nucleus, "H1" and "H2" (atomic numbers $Z_{H1}=1$ and $Z_{H2}=1$) for the hydrogen nuclei.

$$\hat{H}_{H_2O} = -\frac{1}{2M_O}\nabla_{\mathbf{R}_O}^2 - \frac{1}{2M_{H1}}\nabla_{\mathbf{R}_{H1}}^2 - \frac{1}{2M_{H2}}\nabla_{\mathbf{R}_{H2}}^2 - \sum_{i=1}^{10} \frac{1}{2}\nabla_{\mathbf{r}_i}^2$$

Kinetic energy of O
 Kinetic energy of H1
 Kinetic energy of H2
 Kinetic energy of electron i

$$- \sum_{i=1}^{10} \frac{8}{|\mathbf{r}_i - \mathbf{R}_O|} - \sum_{i=1}^{10} \frac{1}{|\mathbf{r}_i - \mathbf{R}_{H1}|} - \sum_{i=1}^{10} \frac{1}{|\mathbf{r}_i - \mathbf{R}_{H2}|}$$

Electron attraction to O
 Electron attraction to H1
 Electron attraction to H2

$$+ \sum_{i=1}^{10} \sum_{j=i+1}^{10} \frac{1}{|\mathbf{r}_i - \mathbf{r}_j|} + \sum_{I=1}^3 \sum_{J=I+1}^3 \frac{Z_I Z_J}{|\mathbf{R}_I - \mathbf{R}_J|}$$

Electron-electron repulsion
 nucleus-nucleus repulsion

E.g. Assume

- $Z_1 = Z_O$
- $Z_2 = Z_{H1}$
- $Z_3 = Z_{H2}$

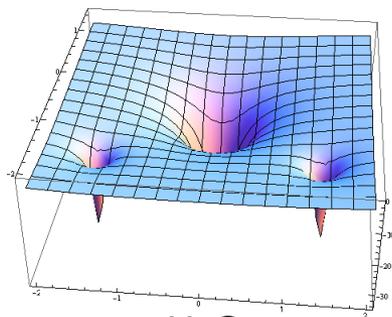
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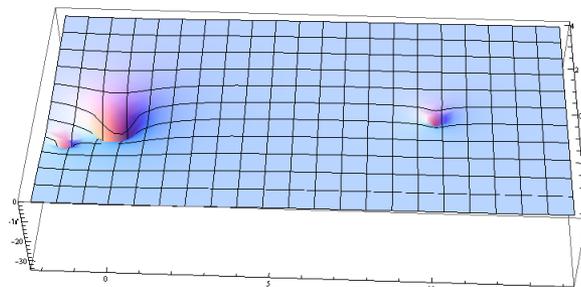
$$\begin{aligned}
 \hat{H}_{\text{H}_2\text{O}} = & \cancel{\frac{1}{2M_O} \nabla_{\mathbf{R}_O}^2}^{\text{B.O.}} - \cancel{\frac{1}{2M_{H1}} \nabla_{\mathbf{R}_{H1}}^2}^{\text{B.O.}} - \cancel{\frac{1}{2M_{H2}} \nabla_{\mathbf{R}_{H2}}^2}^{\text{B.O.}} - \sum_{i=1}^{10} \frac{1}{2} \nabla_{\mathbf{r}_i}^2 \\
 & \text{Kinetic energy of O} \quad \text{Kinetic energy of H1} \quad \text{Kinetic energy of H2} \quad \text{Kinetic energy of electron } i \\
 & - \sum_{i=1}^{10} \frac{8}{|\mathbf{r}_i - \mathbf{R}_O|} - \sum_{i=1}^{10} \frac{1}{|\mathbf{r}_i - \mathbf{R}_{H1}|} - \sum_{i=1}^{10} \frac{1}{|\mathbf{r}_i - \mathbf{R}_{H2}|} \\
 & \text{Electron attraction to O} \quad \text{Electron attraction to H1} \quad \text{Electron attraction to H2} \\
 & + \sum_{i=1}^{10} \sum_{j=i+1}^{10} \frac{1}{|\mathbf{r}_i - \mathbf{r}_j|} + \sum_{I=1}^3 \sum_{J=I+1}^3 \frac{Z_I Z_J}{|\mathbf{R}_I - \mathbf{R}_J|} \\
 & \text{Electron-electron repulsion} \quad \text{nucleus-nucleus repulsion}
 \end{aligned}$$

E.g. Assume
 $Z_1=Z_O$
 $Z_2=Z_{H1}$
 $Z_3=Z_{H2}$

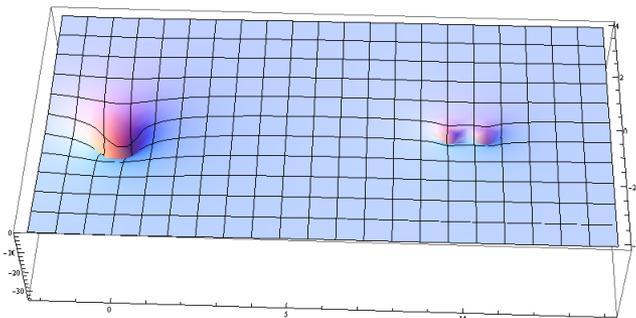
Example: Nuclear attraction potential for one O and two H



H_2O



$\text{OH} + \text{H}$



$\text{O} + \text{H}_2$

- How does the Hamiltonian operator differ between these examples?
- Can you suggest how you may model the reaction $\text{OH} + \text{H} \rightarrow \text{H}_2\text{O}$

Born-Oppenheimer (or adiabatic) approximation

Solve Schrödinger's equation only in the electronic coordinates for each set of given (fixed) nuclear coordinates.

$$\hat{H} = -\sum_I \frac{1}{2M_I} \nabla_I^2 + \sum_I \sum_{J>I} \frac{Z_I Z_J}{|\mathbf{R}_I - \mathbf{R}_J|} - \underbrace{\sum_i \frac{1}{2} \nabla_i^2 + \sum_i \sum_{j>i} \frac{1}{|\mathbf{r}_i - \mathbf{r}_j|} + \sum_i \sum_I -\frac{Z_I}{|\mathbf{r}_i - \mathbf{R}_I|}}_{\hat{H}_{elec}}$$

$$\hat{H}_{elec} \Phi_{elec} = E_{elec} \Phi_{elec}$$

$$\Phi_{elec} = \Phi_{elec}(\{\mathbf{r}_i\}; \{\mathbf{R}_I\})$$

$$E_{elec} = E_{elec}(\{\mathbf{R}_I\})$$

The resulting wavefunctions and energies have a **parametric dependence** in the coordinates of the nuclei

- Central to Chemistry: allows to find reactants, products, transition states, reaction paths
- We will use the Born-Oppenheimer approximation throughout this course

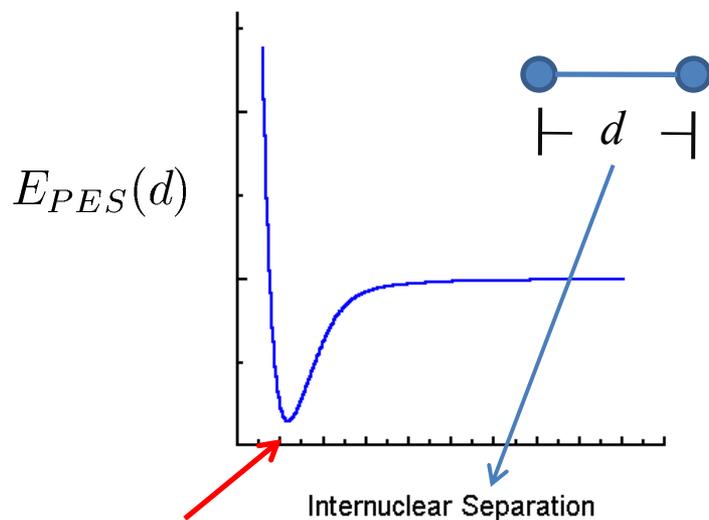
Potential Energy Surface (PES)

The electronic energy is a function of the nuclear positions. So is the internuclear repulsion energy. Their sum is the “total energy” of a molecule under the B.O. approximation. One can represent this as a surface, which is called the potential energy surface (PES).

$$E_{PES}(\{\mathbf{R}_I\}) = E_{elec}(\{\mathbf{R}_I\}) + \sum_I \sum_{J>I} \frac{Z_I Z_J}{|\mathbf{R}_I - \mathbf{R}_J|}$$

Electronic energy Repulsion energy between nuclei I and J

Example: PES of a diatomic

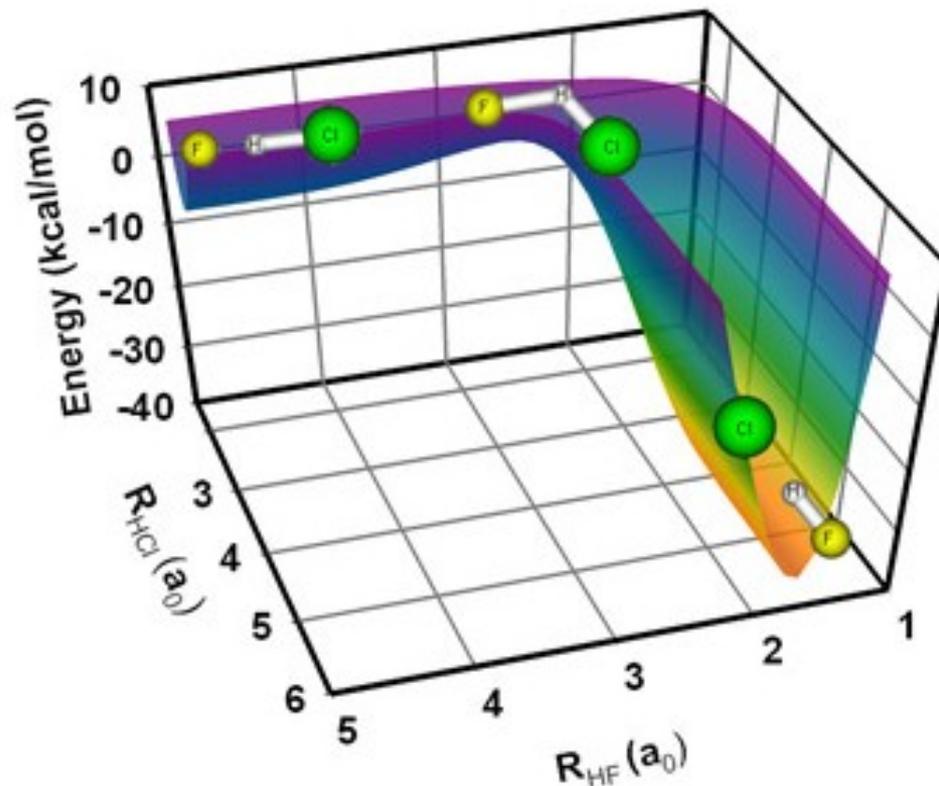


Equilibrium geometry d_e

Reactants, products, transition states, reaction paths

Chemical reactions can be represented as paths followed on a PES

- Reactants and products are minimum energy points on the PES
- Transition states are points on the PES where the energy is maximum along one direction and minimum along all other directions
- We can locate reactants, products and transition states by following paths on the PES until a minimum (or maximum in one direction for TS) is reached (zero partial derivative of energy with respect to each atomic coordinate)



Spin, antisymmetry and the Pauli principle

- Each electron has a “spin”, an intrinsic property which has the characteristics of rotational motion, and is quantised
- Electronic spin is described by a spin angular momentum quantum number $s=1/2$, and its z-component $m_s = 1/2$ (“up” spin) or $-1/2$ (“down” spin)
- We represent the two spin states of the electron by two spin wavefunctions $\alpha(\omega)$ and $\beta(\omega)$ which are orthonormal:

$$\int \alpha^*(\omega)\alpha(\omega)d\omega = \int \beta^*(\omega)\beta(\omega)d\omega = 1$$

$$\langle\alpha|\alpha\rangle = \langle\beta|\beta\rangle = 1$$

$$\int \alpha^*(\omega)\beta(\omega)d\omega = \int \beta^*(\omega)\alpha(\omega)d\omega = 0$$

$$\langle\alpha|\beta\rangle = \langle\beta|\alpha\rangle = 0$$

Spatial and spin coordinates

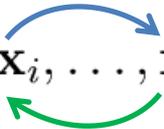
- To describe the spin of an electron we therefore include an extra “spin-coordinate”
- Therefore, an electron is described not only by the three spatial coordinates x,y,z ($=\mathbf{r}$) but also by its spin coordinate ω
- We will denote these four coordinates collectively by \mathbf{x}

$$\mathbf{x} = \{\mathbf{r}, \omega\}$$

- Particles whose spin quantum number s is a half-integer (e.g. $1/2$, $3/2$, $5/2$, etc) are called Fermions
- Electrons and protons are examples of Fermions
- Particles whose spin quantum number s is an integer (e.g. 1 , 2 , 3) are called Bosons. Photons are examples of Bosons

Antisymmetry of electronic wavefunctions

- Wavefunctions of Fermions change sign when the coordinates (space and spin) of any two particles are exchanged.
- This property is called antisymmetry:

$$\Psi(\mathbf{x}_1, \dots, \mathbf{x}_i, \dots, \mathbf{x}_j, \dots, \mathbf{x}_N) = -\Psi(\mathbf{x}_1, \dots, \mathbf{x}_j, \dots, \mathbf{x}_i, \dots, \mathbf{x}_N)$$


- Electrons are Fermions, therefore electronic wavefunctions **must be antisymmetric**
- We need to include antisymmetry in all approximate wavefunctions we construct

Summary / Reading assignment

- Born-Oppenheimer approximation (Cramer, page 110)
- Many-electron wavefunctions (Cramer, pages 119-122)
- Antisymmetry of electronic wavefunctions (Cramer, pages 122-126)