

# Hyperbolic mass and Maskit gluings

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Southampton, March 2022

based on joint work with Erwann Delay and Rafaela Wutte

# Energy for **locally asymptotically hyperbolic** manifolds

space-dimension  $n$

Theorem (with E. Delay, arXiv:1901.05263)

*The energy-momentum vector of conformally compact  $n$ -dimensional asymptotically locally hyperbolic manifolds  $(M, g)$  with **spherical infinity** and with scalar curvature  $R(g)$  satisfying  $R(g) \geq -n(n-1)$ ,  $n \geq 3$ , is **timelike future-pointing or vanishes**.*

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Theorem (with E. Delay and R. Wutte, arxiv:2112.00095)

*There exist 3-dimensional conformally compact asymptotically locally hyperbolic Riemannian manifolds  $(M, g)$  with*

$$R(g) = -6,$$

*with arbitrarily high genus and with **negative total mass***

# Positive energy for asymptotically hyperbolic manifolds

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- 2 Different story if topology at infinity is not spherical.
- 3 Huang, Jang, Martin (2019): lightlike cannot occur
- 4 Generalises to many ends and boundaries with  $H < n-1$  (PTC, Galloway, 2107.05603 [gr-qc])

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- 4 Generalises to many ends and boundaries with  $H < n - 1$  (PTC, Galloway, 2107.05603 [gr-qc])
- 5 if  $n \geq 7$ , needs the higher-dimensional asymptotically flat positive energy theorem (Lohkamp, Schoen & Yau)

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- 5 key idea: the “Maskit gluing” by Isenberg, Lee & Stavrov (2010)

# Other topologies at infinity?

In dimension  $3 + 1$ :

- positivity of energy in the spherical case is taken care of by the last theorem
- toroidal case: the Horowitz-Myers metrics provide examples with **negative mass** without black hole boundaries
- this talk: **what about higher genus at infinity** without black hole boundaries?

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time-symmetric vacuum general relativistic initial data with suitably normalised negative cosmological constant

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previously: toroidal infinity

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not clear how to generalise this to higher dims

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Why should we care?

- *Asymptotically hyperbolic manifolds* are ubiquitous in nowadays theoretical physics (supergravities, string theory, holography, CFT/AdS).
- They appear naturally as spacelike hypersurfaces in solutions of Einstein equations, with or without a cosmological constant  $\Lambda$ :  
hyperbolic space itself occurs as a “*static slice*” of the Anti-de Sitter spacetime ( $\Lambda < 0$ ), or as a *hyperboloid* in Minkowski spacetime  $\Lambda = 0$ .
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# Model metrics: Kottler-Birmingham metrics

Static vacuum solutions of Einstein equations with a negative cosmological constant

$$\mathbf{g}_m = -V_m^2 dt^2 + V_m^{-2} dr^2 + r^2 h_{\kappa}, \quad V_m^2 = r^2 + \kappa - \frac{2m}{r^{n-2}}.$$

where  $h_{\kappa}$  is a  $t$ - and  $r$ -independent Einstein metric on a  $(n-1)$ -dim compact manifold, with scalar curvature  $R(h) = (n-1)(n-2)\kappa$ .

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- *asymptotically BK* is the same as **locally asymptotically hyperbolic** in space-dimension 3
- and is a special case of **locally asymptotically hyperbolic** in higher dimensions



# Model metrics: Horowitz-Myers Instantons

$$\mathbf{g}_m = -V_m^2 dt^2 + V_m^{-2} dr^2 + r^2(d\theta^2 + h'_0), \quad V_m^2 = r^2 \cancel{A/k} - \frac{2m}{r^{n-2}}.$$

where  $h'_0$  is a  $t$ -,  $\theta$ -, and  $r$ -independent **Ricci flat** metric on a  $(n - 3)$ -dim compact manifold.



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Woolgar's version of the Horowitz-Myers conjecture

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- The mass relative to  $g_0$  can be arbitrarily negative, proportional to the negative of  $m$ .
- Horowitz-Myers conjecture: these are minima of energy at prescribed conformal structure at infinity.

# Idea: use “gluing at infinity”

“Maskit gluing”

Theorem (Isenberg, Lee & Stavrov 2010, PTC, Delay, arXiv:1511.07858)

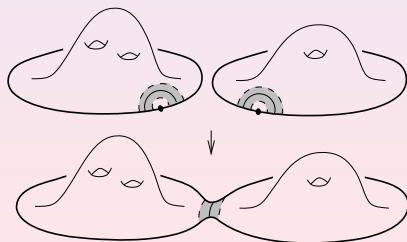
*Given two asymptotically hyperbolic manifolds with constant scalar curvature (or general relativistic vacuum initial data sets) one can construct a new one by making a connected sum at the conformal boundary at infinity. The construction can be localised by a Carlotto-Schoen type hyperbolic gluing.*

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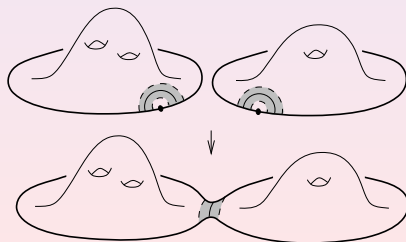


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**Question:** What is the energy-momentum of the new initial data set?

# How to define mass

## Spacetime methods

- 1 Spacetime variational methods: “Noether charge” *à la Wald* ( $\sim 1990$ )  $\equiv$  geometric Hamiltonian methods *à la Kijowski-Tulczyjew* (1979)

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- 1 Spacetime variational methods: “Noether charge” *à la Wald* ( $\sim 1990$ )  $\equiv$  geometric Hamiltonian methods *à la Kijowski-Tulczyjew* (1979)
- 2 A convenient geometric formula for total energy  $E$ :  
if  $g$  approaches a *Kottler-Birmingham* metric with  $m = 0$

$$E = -\frac{1}{16(n-2)\pi} \lim_{R \rightarrow \infty} \int_{r=R} D^j V (\mathbf{R}^i_j - \frac{\mathbf{R}}{n} \delta^i_j) dS_i.$$

where  $\mathbf{R}^i_j$  is the Ricci tensor of  $g$  and

$$V = \sqrt{r^2 + \kappa}, \quad \kappa \in \{0, \pm 1\}. \quad (**)$$

# Energy-momentum, spherical conformal infinity

Total energy-momentum  $p_{(\mu)}$ :

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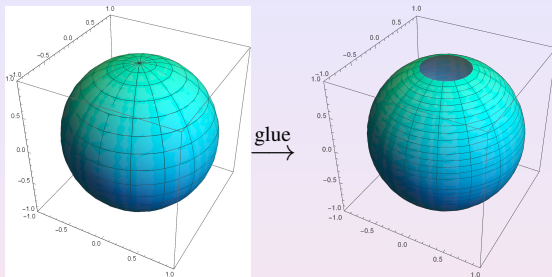
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- $p_{(\mu)}$  transforms as a Lorentz vector under isometries of the hyperbolic metric.

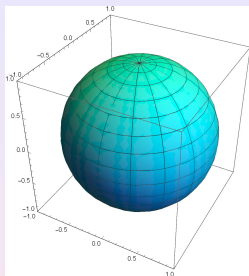
# Positive energy for asymptotically hyperbolic manifolds

Energy-momentum vector and localised Maskit gluing

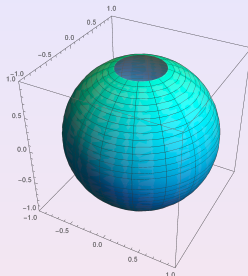


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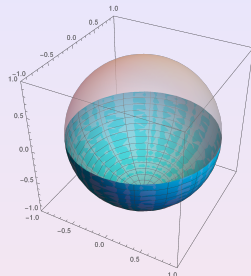
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glue



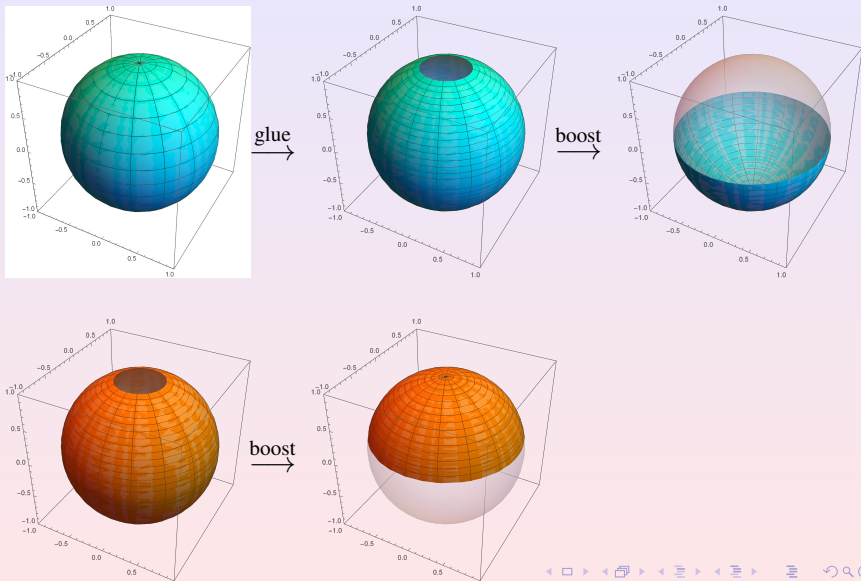
boost





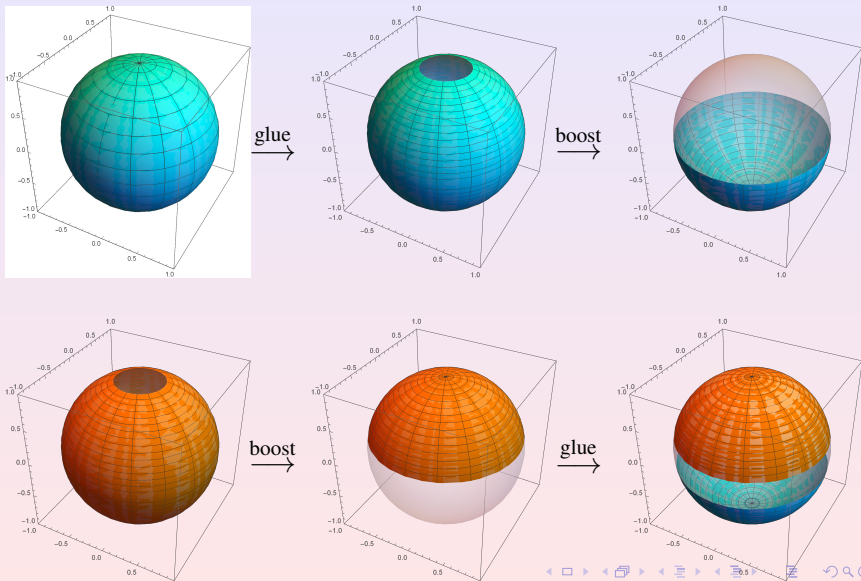
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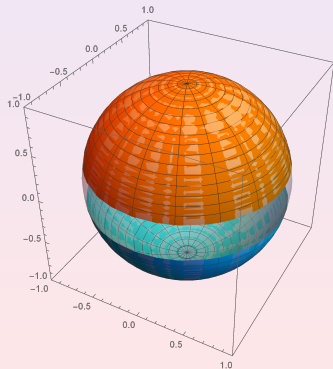


# Now energy-momentum is obviously additive

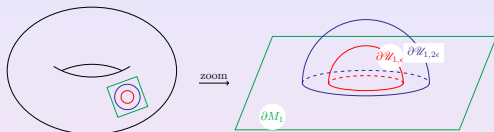
$$p_{(\mu)} = -\frac{1}{16(n-2)\pi} \lim_{R \rightarrow \infty} \int_{r=R} D^j V_{(\mu)} \left( R^i_j - \frac{R}{n} \delta^i_j \right) dS_i.$$

where

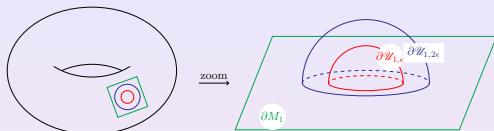
$$V_{(0)} = \sqrt{r^2 + 1}, \quad V_{(i)} = x^i.$$



# Carlotto-Schoen type gluing, toroidal infinity

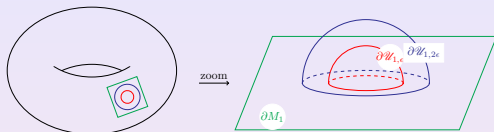


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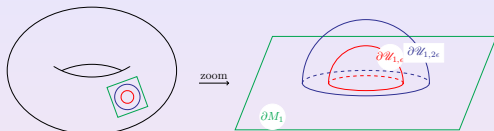
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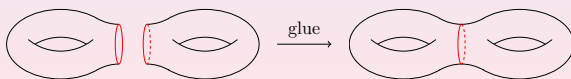


- the metric is exactly hyperbolic inside the red half-ball
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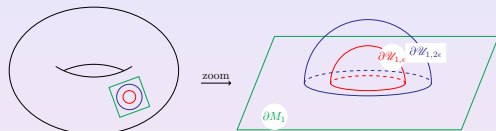
# Carlotto-Schoen type gluing, toroidal infinity



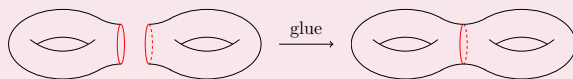
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# Carlotto-Schoen type gluing, toroidal infinity



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- the boundary of the red half-ball is totally geodesic
- the hyperbolic metric extends smoothly when any two such boundaries touch



- the *initial mass* is defined with respect to a **toroidal** BK metric; the *final one* with respect to a **genus-two** BK metric!



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$$E = -\frac{1}{16\pi} \lim_{R \rightarrow \infty} \int_{r=R} D^j r \left( R^i_j - \frac{R}{3} \delta^i_j \right) dS_i.$$

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$$\begin{aligned} E &= -\frac{1}{16\pi} \lim_{R \rightarrow \infty} \int_{\bar{r}=R} D^j(\sqrt{\bar{r}^2 - 1}) (\mathbf{R}^i_j - \frac{R}{3} \delta^i_j) dS_i \\ &= -\frac{1}{16\pi} \lim_{R \rightarrow \infty} \int_{r=R} D^j(e^{-\omega/2} r) (\mathbf{R}^i_j - \frac{R}{3} \delta^i_j) dS_i. \end{aligned}$$



# Mass formula, space dimensions 3, somewhat more generally:

## Theorem

Let  $g$  be *asymptotic to two backgrounds*,

$$b = \frac{dr^2}{r^2 + \kappa} + r^2 h_\kappa \text{ and } \bar{b} = \frac{d\bar{r}^2}{\bar{r}^2 + \bar{\kappa}} + \bar{r}^2 h_{\bar{\kappa}}, \text{ with } h_{\bar{\kappa}} = e^\omega h_\kappa.$$

Then

$$E = -\frac{1}{16\pi} \lim_{R \rightarrow \infty} \int_{r=R} D^j ({}^R r^i_j - \frac{R}{3} \delta^i_j) dS_i.$$

$$\bar{E} = -\frac{1}{16\pi} \lim_{R \rightarrow \infty} \int_{r=R} D^j (e^{-\omega/2} r^i_j) ({}^R r^i_j - \frac{R}{3} \delta^i_j) dS_i.$$

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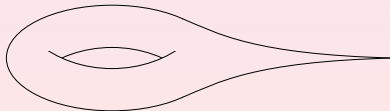
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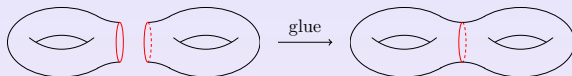
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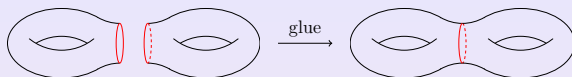
# Gluings torii



- mass of each half of the glued manifold

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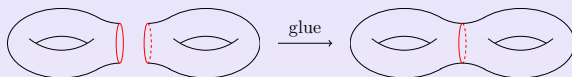
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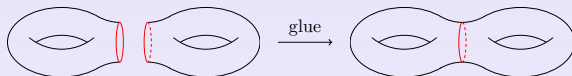
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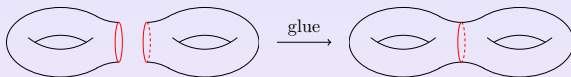
# Gluing torii, $\epsilon$ small needed



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# Gluing torii, $\text{limit } \epsilon \rightarrow 0$ needed

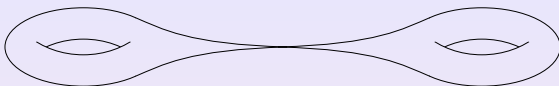


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# Taking the limit $\epsilon \rightarrow 0$

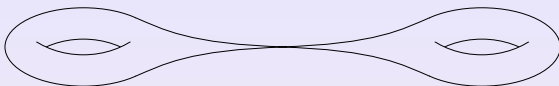


## Theorem (PTC, E. Delay, R. Wutte)

When *Maskit-gluing* two Horowitz-Myers metrics with mass parameter  $m$ ,  $e^{\omega_\epsilon}$  tends to the conformal factor  $e^{\omega_0}$  of a punctured torus as  $\epsilon$  tends to zero, with

$$\bar{E} \rightarrow -\frac{m}{4\pi} \int_{\mathbb{T}^2} e^{-\omega_0/2} d\mu_{h_0} < 0 \quad (1)$$

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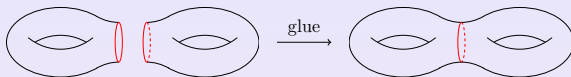
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It thus follows that the final mass is negative for  $\epsilon$  small enough

# Taking the limit $\epsilon \rightarrow 0$



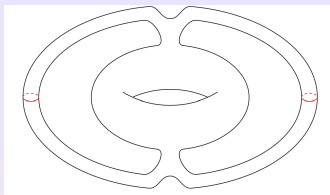
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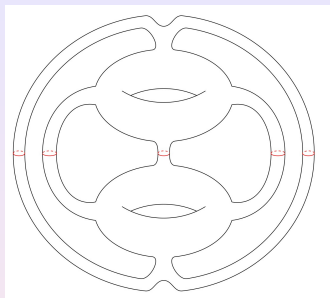
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The construction can be iterated

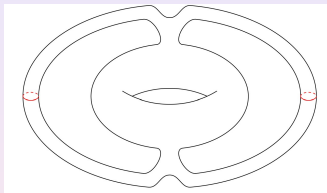
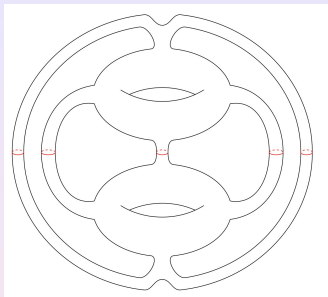
# Gluings with several punctures



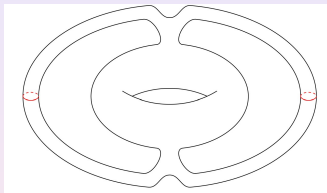
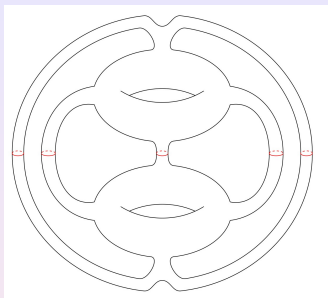
# Gluings with several punctures



# “Topological instability at the conformal boundary”?

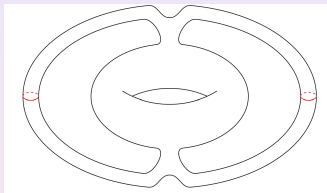
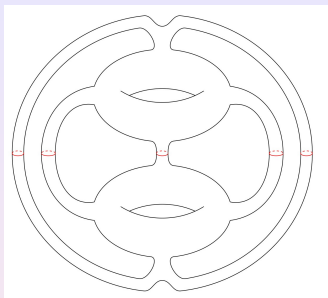


# “Topological instability at the conformal boundary”?



The above construction can be used to **lower** the total mass of an ALH manifold by a localised deformation near the conformal boundary at infinity, *for geometries with very thin necks*

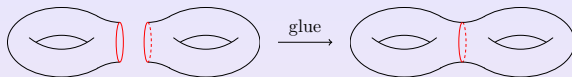
# “Topological instability at the conformal boundary”?



The **existing higher-genus-mass-inequalities**, which include conditions such as *existence of a strictly negative mass aspect function* (Lee & Neves; Gibbons), or *product topology* (Galloway et al.), **cannot be improved without further conditions**



# Conjectures:

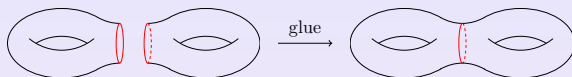


- For any genus of the conformal boundary at infinity there exists  $m_c \leq 0$ , depending only upon the conformal class of conformal infinity, so that for conformally compact vacuum asymptotically locally hyperbolic initial data sets we have

$$E \geq m_c,$$

with  $m_c < 0$  unless the boundary is spherical.

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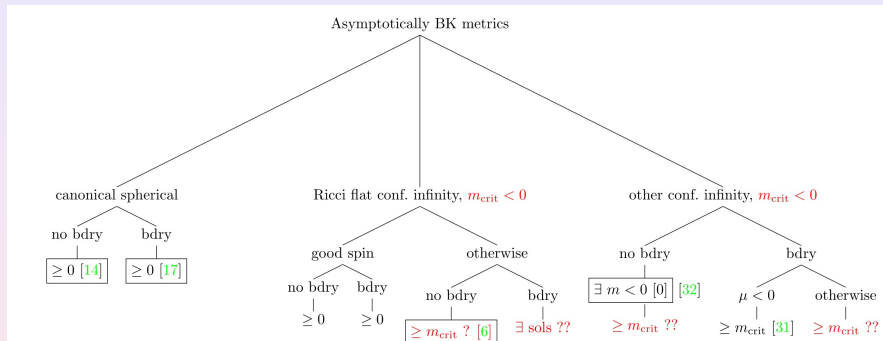
- $m_c$  is *attained on a static metric*.

# Hyperbolic mass, asymptotically Birmingham-Kottler metrics

Conformally compact, with or without black-hole boundary

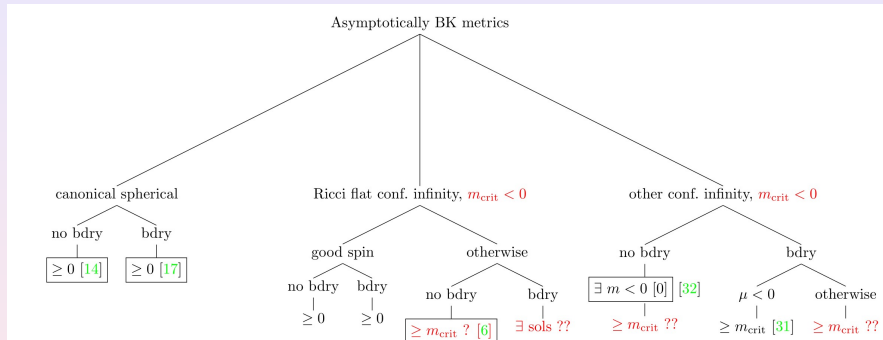
# Hyperbolic mass, asymptotically Birkhoff-Kottler metrics

Conformally compact, with or without black-hole boundary



# Hyperbolic mass, asymptotically Birminghams-Kottler metrics

Conformally compact, with or without black-hole boundary

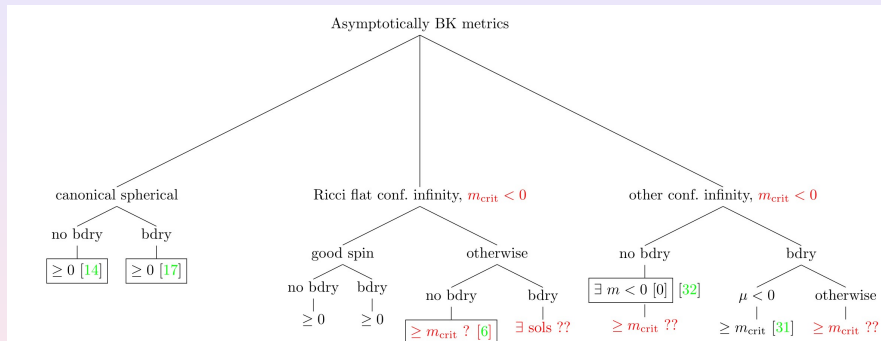


Negative mass solutions:

- toroidal: Horowitz-Myers (1998)

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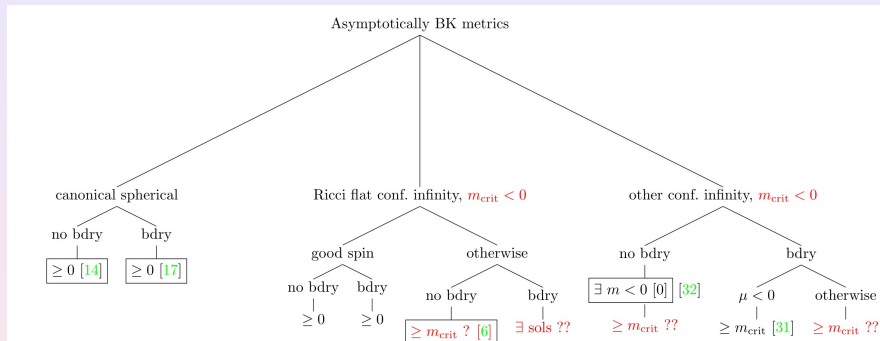


Negative mass solutions:

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Conformally compact, with or without black-hole boundary



Negative mass solutions:

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- quotients of a sphere: Clarkson & Mann (2006), dim 4+1
- higher genus: PTC, Delay, Wutte (XII 2021), dim 3+1