Fermionic spinfoam models and TQFTs

Steven Kerr

University of Nottingham

Sum over Histories

Sum over Histories

$$\mathbb{Z}=\text{``}\int\mathcal{D}g\;e^{iS}$$
''

Sum over Histories

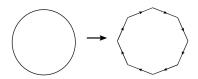
$$\mathbb{Z}=\text{``}\int\mathcal{D}g\;e^{iS\text{''}}$$

Fundamental length scale

Sum over Histories

$$\mathbb{Z}=\text{``}\int\mathcal{D}g\;e^{iS\text{''}}$$

► Fundamental length scale



Problems

► Triangulation independence

Problems

- ► Triangulation independence
- Absence of matter

Problems

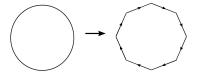
- Triangulation independence
- Absence of matter
- ► We take the point of view that matter and triangulation independence are crucial!

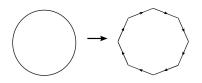
$$\mathbb{Z} = \int \mathcal{D}\psi \mathcal{D}\overline{\psi} \ \mathrm{e}^{i\int \, \overline{\psi} \not \! D\psi} \qquad \not \! D = \gamma^\mu (d_\mu - i A_\mu)$$

$$egin{aligned} \mathbb{Z} &= \int \mathcal{D}\psi \mathcal{D}\overline{\psi} \ e^{i\int \, \overline{\psi} \not\!\!\!\!/ \psi} & \not\!\!\!\!/ = \gamma^\mu (d_\mu - i A_\mu) \ &= \det(i \not\!\!\!/ D) = e^{\mathrm{tr} \ln i \not\!\!\!/ D} \end{aligned}$$

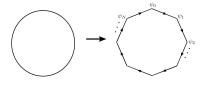
$$egin{aligned} \mathbb{Z} &= \int \mathcal{D}\psi \mathcal{D}\overline{\psi} \, e^{i\int \, \overline{\psi} \not\!\!\!/ \psi} & \not\!\!\!/ egin{aligned} eta &= \gamma^{\mu} (d_{\mu} - i \!\!\!/ A_{\mu}) \ &= \det(i \not\!\!\!/ D) = e^{\mathrm{t} r \ln i \not\!\!\!/ D} \ &= e^{i \!\!\!/ S_{\mathrm{eff}}} \end{aligned}$$

▶ John Barret has suggested that the Standard Model can be induced in this way: arXiv:1101.6078v2 [hep-th]

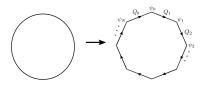




$$\psi(t), \bar{\psi}(t), t \in [0, 2\pi]$$
 $\psi_i, \bar{\psi}_i, i = 1..N$



$$\psi(t), \bar{\psi}(t), t \in [0, 2\pi]$$
 $\psi_i, \bar{\psi}_i, i = 1..N$



$$\psi(t), \bar{\psi}(t), t \in [0, 2\pi]$$
 $\psi_i, \bar{\psi}_i, i = 1..N$

$$A(t)$$
 $Q_i = \mathcal{P}e^{i\int Adt}$

$$egin{aligned} \hat{\mathbb{Z}} &= \int \prod_{i=1}^{N} d\psi_i dar{\psi}_i \; e^{\sum_{i=1}^{N} ar{\psi}_i (\psi_i - Q_{i+1}\psi_{i+1})} \ \psi_{N+1} &= \psi_1 \quad Q_{N+1} = Q_1 \end{aligned}$$

$$egin{aligned} \hat{\mathbb{Z}} &= \int \prod_{i=1}^{N} d\psi_i dar{\psi}_i \; e^{\sum_{i=1}^{N} ar{\psi}_i (\psi_i - Q_{i+1} \psi_{i+1})} \ \psi_{N+1} &= \psi_1 \quad Q_{N+1} = Q_1 \ &= \det(1-Q) \qquad Q = \prod_{i=1}^{N} Q_i \end{aligned}$$

$$egin{aligned} \hat{\mathbb{Z}} &= \int \prod_{i=1}^N d\psi_i dar{\psi}_i \; e^{\sum_{i=1}^N ar{\psi}_i (\psi_i - Q_{i+1}\psi_{i+1})} \ \psi_{N+1} &= \psi_1 \quad Q_{N+1} = Q_1 \ &= \det(1-Q) \qquad \qquad Q = \prod_{i=1}^N Q_i \end{aligned}$$

 $ightharpoonup \hat{\mathbb{Z}}$ is triangulation independent - a topological invariant!

Action

What is the significance of this theory? It is a discretisation of a one dimensional Dirac theory,

$$\hat{S} = -i \sum_{i=1}^{N} \bar{\psi}_i \left(\psi_i - Q_{i+1} \psi_{i+1} \right)$$
$$= i \Delta t \sum_{i=1}^{N} \bar{\psi}_i \left(\frac{Q_{i+1} \psi_{i+1} - \psi_i}{\Delta t} \right) \quad \Delta t = \frac{2\pi}{N}$$

Action

What is the significance of this theory? It is a discretisation of a one dimensional Dirac theory,

$$\hat{S} = -i \sum_{i=1}^{N} \bar{\psi}_{i} \left(\psi_{i} - Q_{i+1} \psi_{i+1} \right)$$

$$= i \Delta t \sum_{i=1}^{N} \bar{\psi}_{i} \left(\frac{Q_{i+1} \psi_{i+1} - \psi_{i}}{\Delta t} \right) \quad \Delta t = \frac{2\pi}{N}$$

$$\lim_{\Delta t \to 0} i \left(\frac{Q_{i+1} \psi_{i+1} - \psi_{i}}{\Delta t} \right) = \not D_{t} \psi(t)$$

$$\lim_{\Delta t \to 0} \Delta t \sum_{i=1}^{N} = \int_{0}^{2\pi} dt$$

$$\lim_{\Delta t \to 0} \hat{S} = \int dt \overline{\psi} \not D \psi$$

One can calculate the partition function of the continuum theory exactly

One can calculate the partition function of the continuum theory exactly

$$\mathbb{Z} = \int \mathcal{D}\psi \mathcal{D}\overline{\psi} \ \mathsf{e}^{\mathsf{i}\int \mathsf{d} \mathsf{t} \, \overline{\psi} \not \! D \psi}$$

One can calculate the partition function of the continuum theory exactly

$$\mathbb{Z} = \int \mathcal{D}\psi \mathcal{D}\overline{\psi} \, \mathsf{e}^{\mathsf{i}\int \mathsf{d}t\,\overline{\psi} \not{\!\!\!D}\psi}$$

We find that $\hat{\mathbb{Z}} = \mathbb{Z}!$

One can calculate the partition function of the continuum theory exactly

$$\mathbb{Z} = \int \mathcal{D}\psi \mathcal{D}\overline{\psi} \, \mathsf{e}^{i\int dt \, \overline{\psi}
ot\!\!/ \psi}$$

We find that $\hat{\mathbb{Z}} = \mathbb{Z}!$

Naturally, one would like to try do something similar in higher dimensions. This is the subject of current investigation.

Thanks!