

# Self-consistent orbital evolution of a particle around a Schwarzschild black hole

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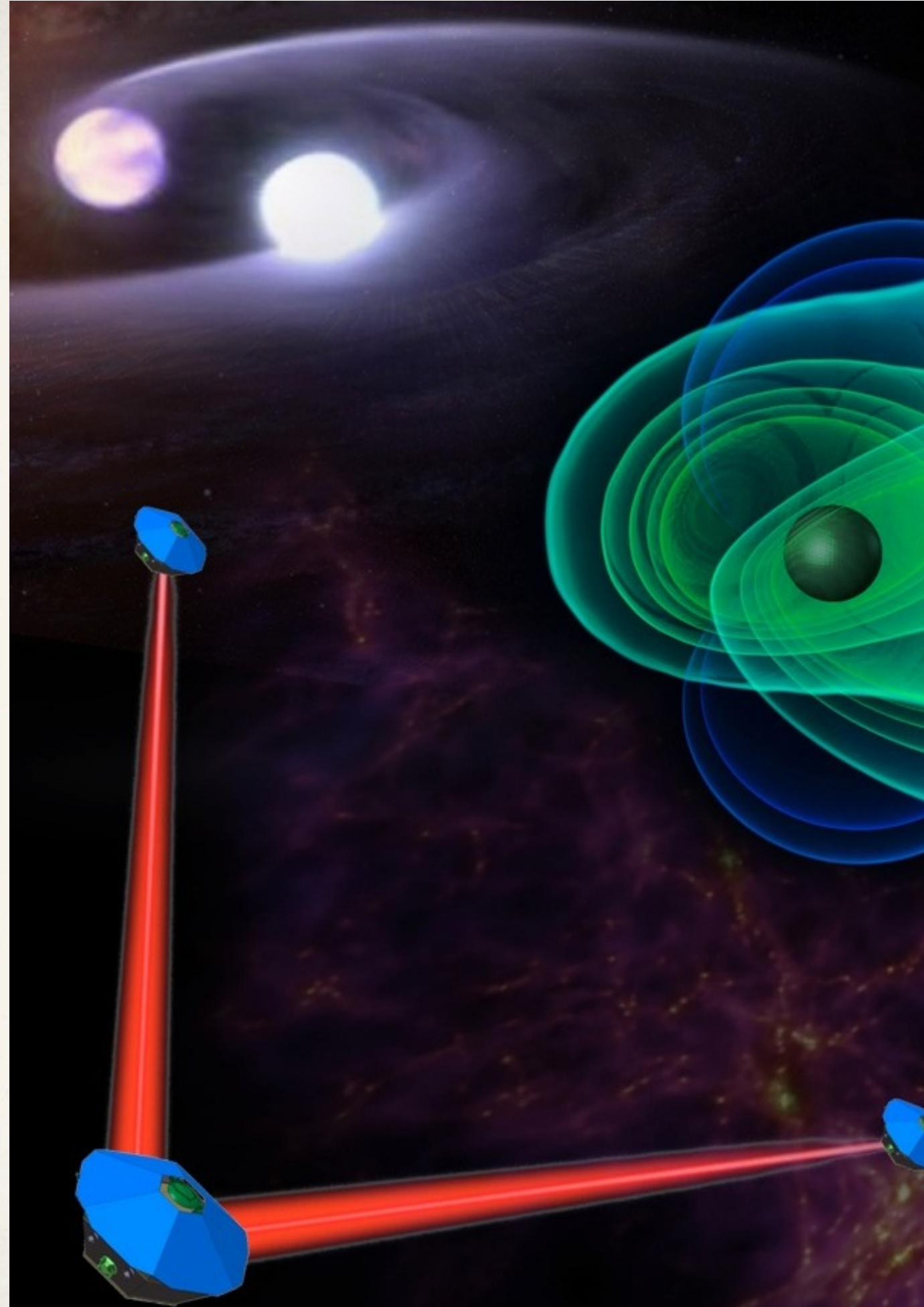
Collaborators: Peter Diener, Ian Vega, Steven Detweiler



# EMRIs and eLISA/NGO

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- ❖ Extreme Mass Ratio Inspirals have long been a promising source of gravitational waves for the LISA, the space based gravitational wave detector.
- ❖ Accurate models are a critical component of any observation.
- ❖ Even more true now that LISA is no more and there are proposals for eLISA/NGO which will have less sensitivity.





# Motion of a point particle

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- ❖ Solve the coupled system of equations for the motion of the particle and its retarded field.
- ❖ Self-interaction of the particle with its retarded field,  $\Phi^{\text{ret}}$ .
- ❖  $\Phi^{\text{ret}}$  diverges like  $1/r$  on the world-line.
- ❖ “Unphysical” divergence removed by appropriate regularization.

$$\square \Phi^{\text{ret}} = -4\pi q \int \frac{\delta^4(x - z(\tau))}{\sqrt{-g}} d\tau$$
$$\frac{Du^\alpha}{d\tau} = a^\alpha = \frac{\bar{q}}{m(\tau)} (g^{\alpha\beta} + u^\alpha u^\beta) \nabla_\beta \Phi^{\text{ret}}$$
$$\frac{dm}{d\tau} = -\bar{q} u^\beta \nabla_\beta \Phi^{\text{ret}}$$



# Effective source regularization

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- ❖ Split retarded field into locally constructed field and “regularized” remainder.
- ❖ Derive an equation for  $\Phi^R$ .
- ❖ Always work with  $\Phi^R$  instead of  $\Phi^{\text{ret}}$ .
- ❖ If  $\Phi^S$  is chosen appropriately, then we can just replace  $\Phi^{\text{ret}}$  with  $\Phi^R$  in the equations of motion.

$$\Phi^{\text{ret}} = \Phi^S + \Phi^R$$

$$\square\Phi^R = \square\Phi^{\text{ret}} - \square\Phi^S$$

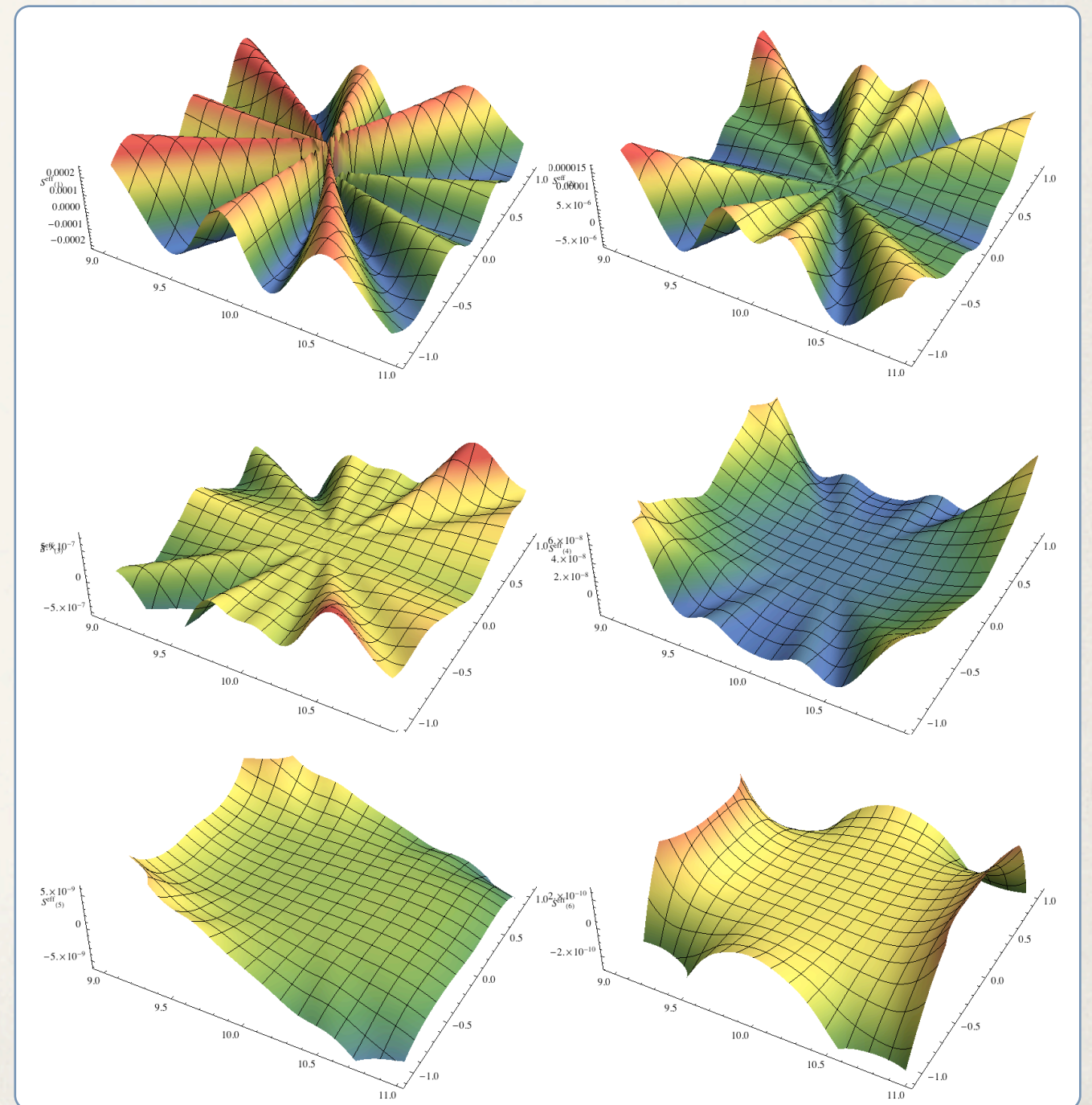
$$\frac{Du^\alpha}{d\tau} = a^\alpha = \frac{\bar{q}}{m(\tau)} (g^{\alpha\beta} + u^\alpha u^\beta) \nabla_\beta \Phi^R$$

$$\frac{dm}{d\tau} = -\bar{q}u^\beta \nabla_\beta \Phi^R$$



# Effective source regularization

- ❖ If  $\Phi^S$  is exactly the Detweiler-Whiting singular field,  $\Phi^R$  is a solution of the homogeneous wave equation.
- ❖ If  $\Phi^S$  is only approximately the Detweiler-Whiting singular field, then the equation for  $\Phi^R$  has an effective source,  $S$ .
- ❖  $S$  is typically finite, but of limited differentiability on the world line.





# Self-consistent Evolution

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- ❖ Solve the coupled system of equations for the motion of the particle and its regularized field.
- ❖  $\Phi^R = \Phi^{\text{ret}}$  in the wave zone
- ❖  $\Phi^R$  finite and (typically) twice differentiable on the world-line

$$\square\Phi^R = S(x|z(\tau), u(\tau))$$

$$\frac{Du^\alpha}{d\tau} = a^\alpha = \frac{\bar{q}}{m(\tau)} (g^{\alpha\beta} + u^\alpha u^\beta) \nabla_\beta \Phi^R$$

$$\frac{dm}{d\tau} = -\bar{q}u^\beta \nabla_\beta \Phi^R$$

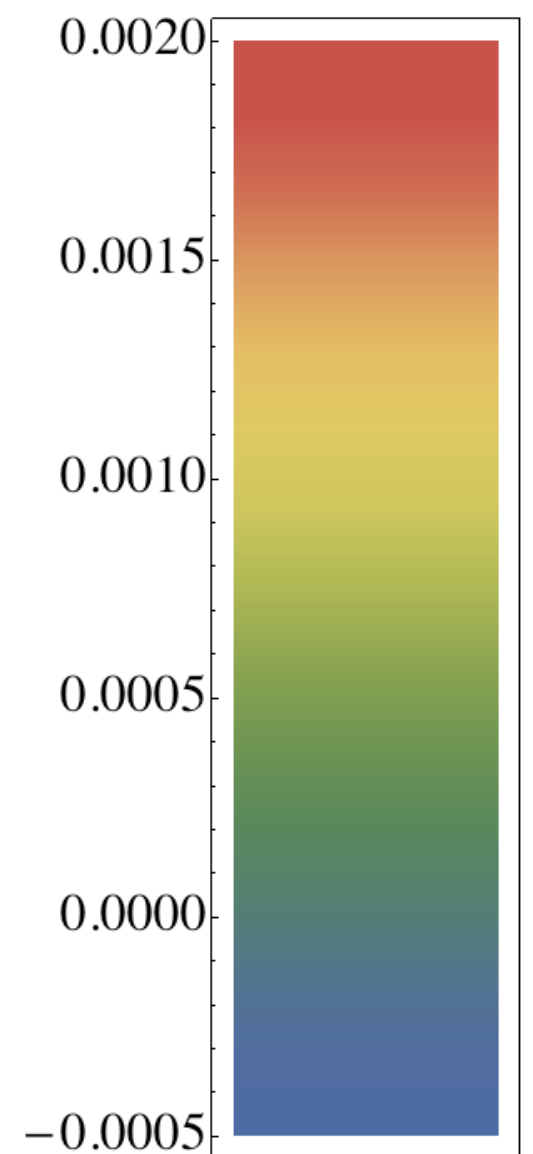
$$\Phi^R(t)$$

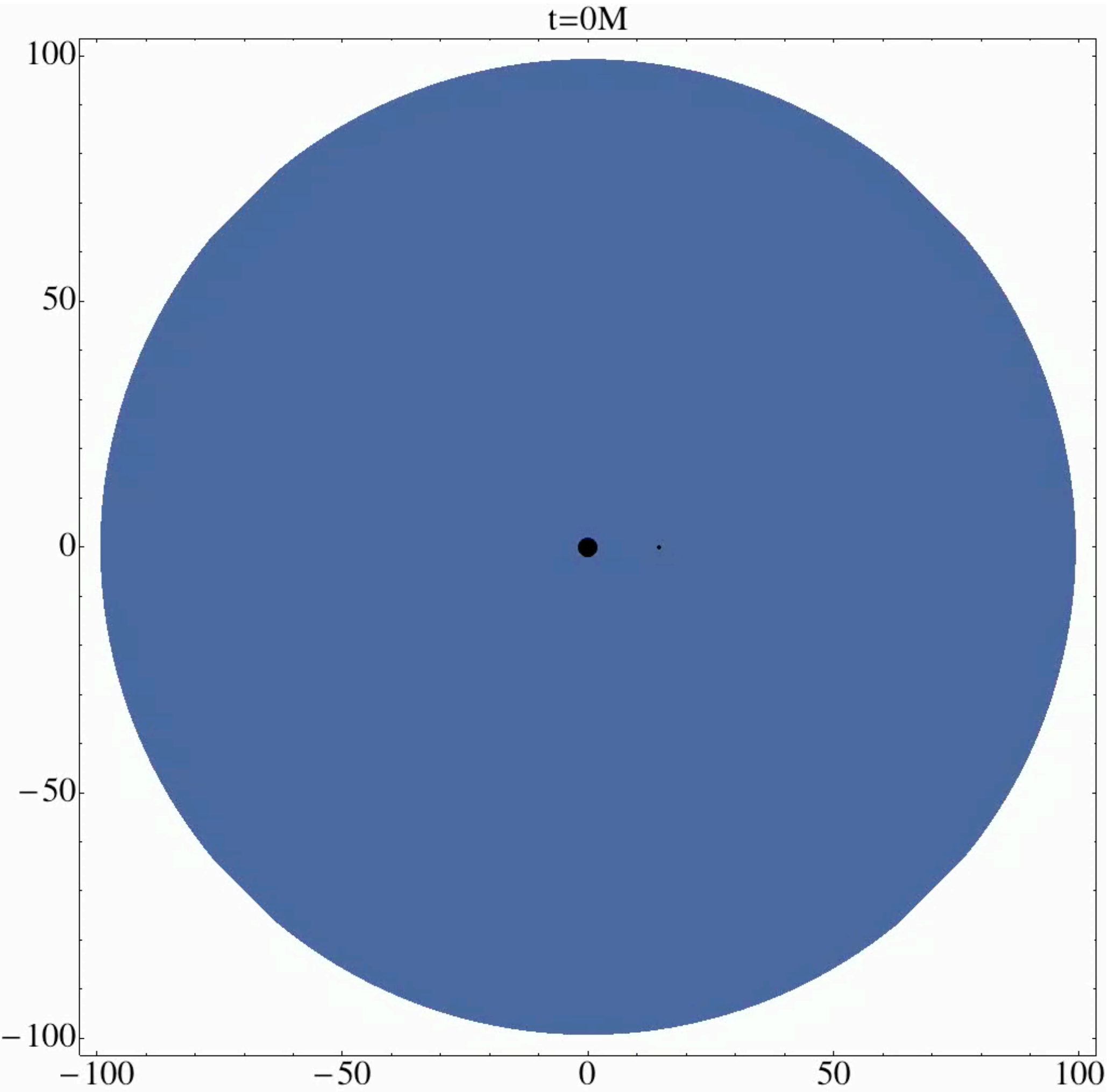
$$q = M/32$$

$$m = M$$

$$p_0 = 7.2$$

$$e_0 = 0.5$$





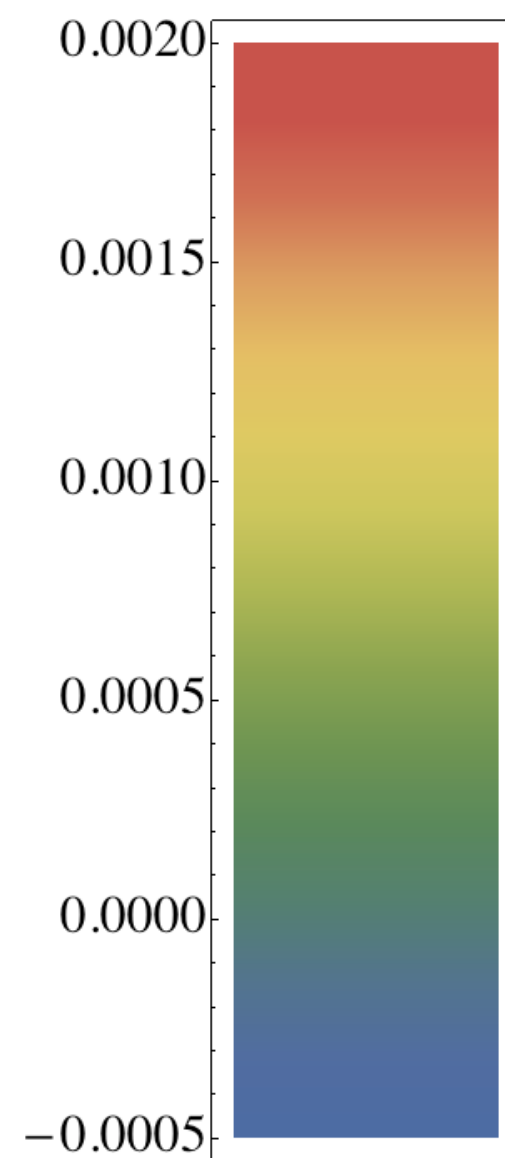
$\Phi^R(t)$

$$q = M/32$$

$$m = M$$

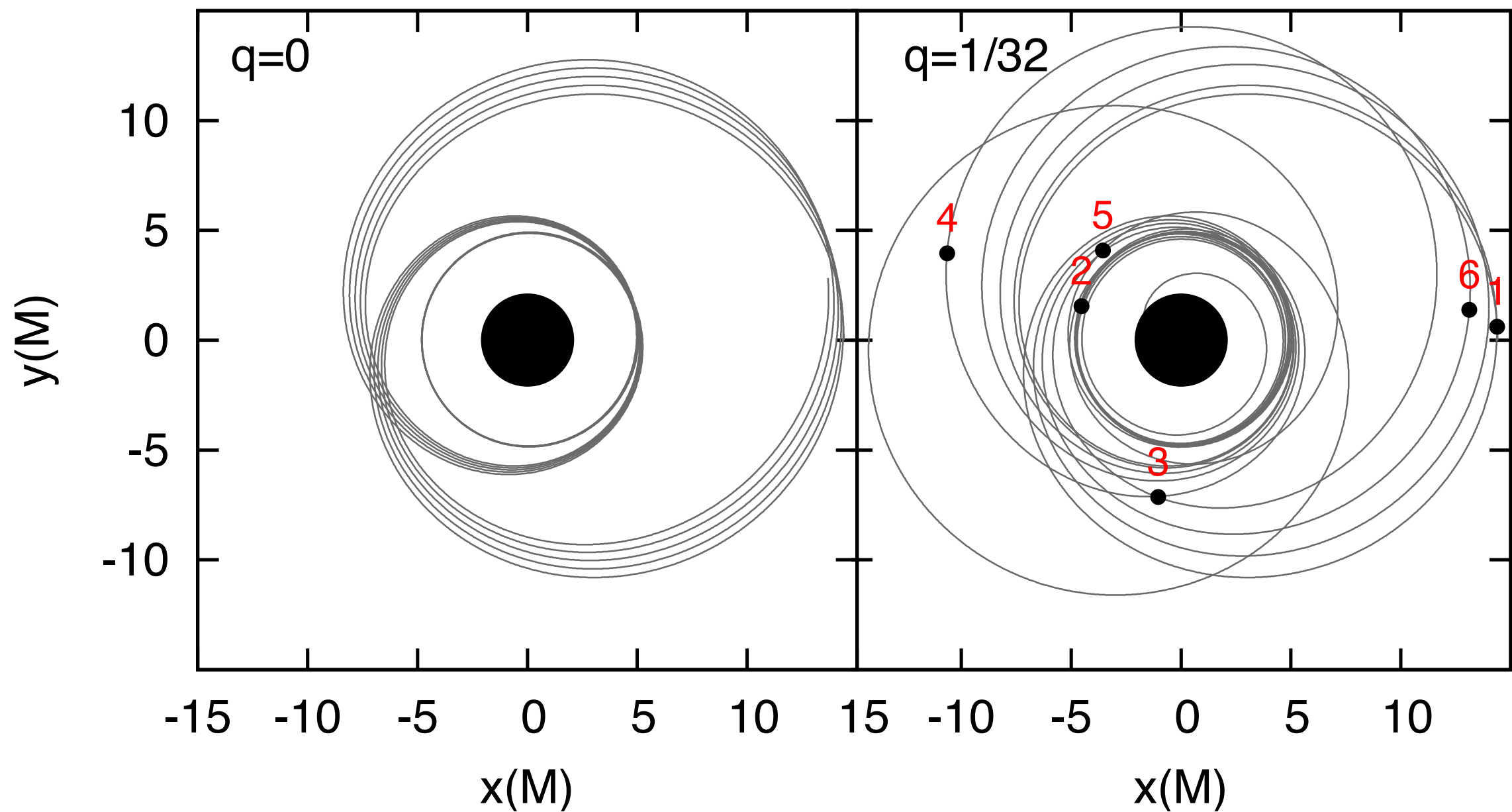
$$p_0 = 7.2$$

$$e_0 = 0.5$$





# Orbital motion



# Orbital motion

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- ❖ Parametrize orbits in terms of a dimensionless semilatus rectum  $p$  and eccentricity  $e$ , such that  $r_{\pm} = Mp / (1 \mp e)$ .

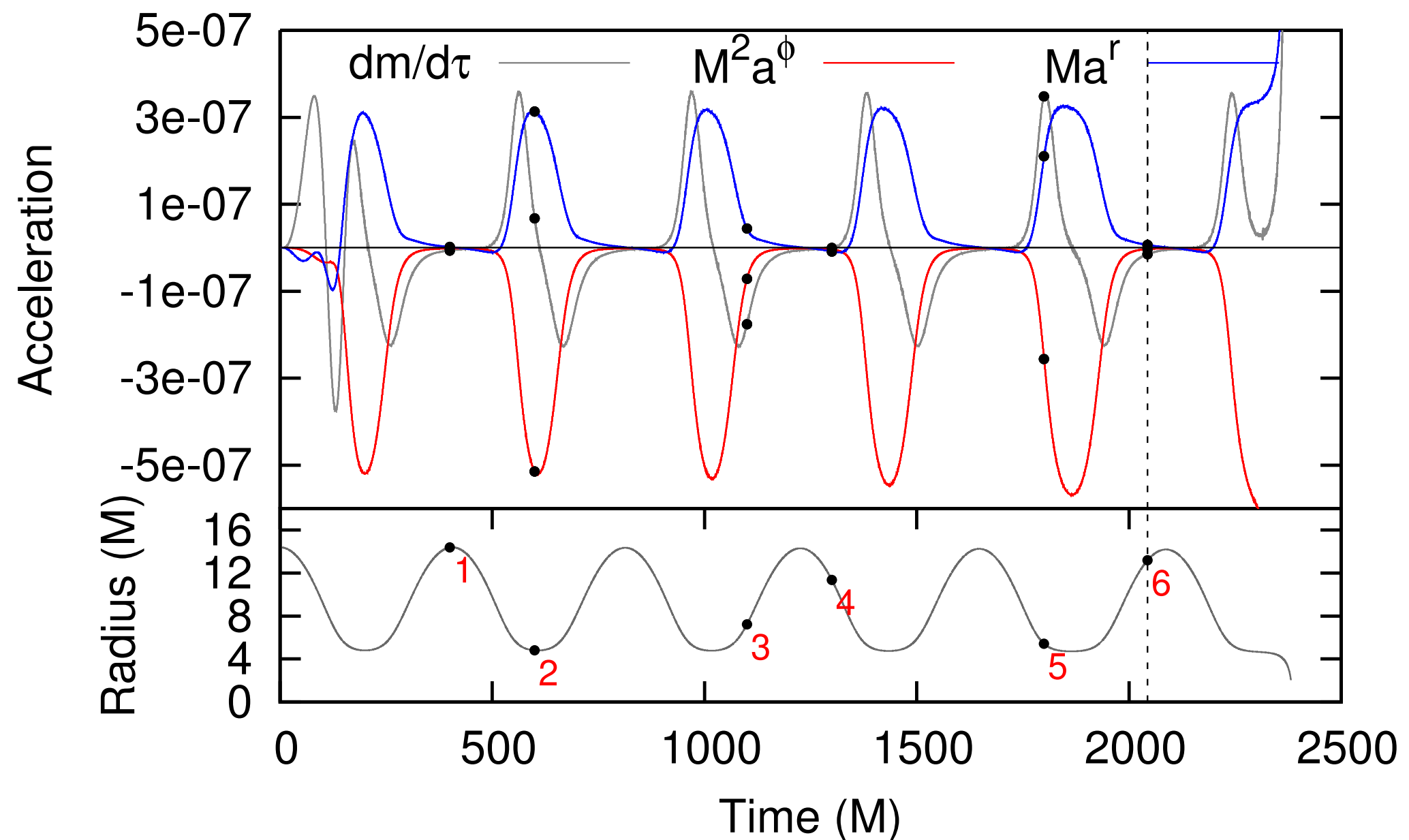
$$r(t) = \frac{Mp}{1 + e \cos(\chi - w)}$$

- ❖ Separatrix,  $p = 6 + 2e$ , corresponds to unstable circular orbits and represents the boundary in  $p$ - $e$  space separating bound from plunging orbits.

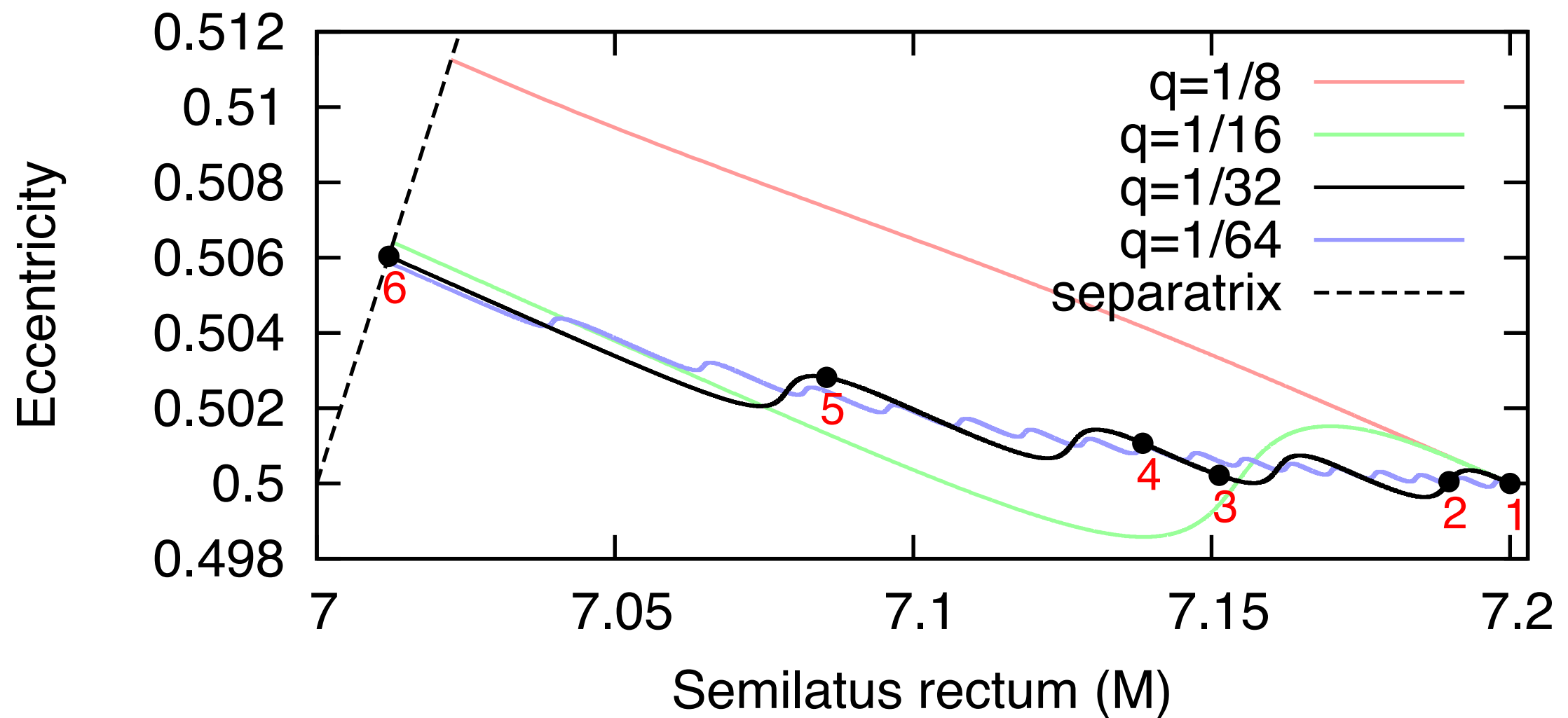
$$\frac{d\phi}{dt} = \left[ 1 - \frac{2Mr'}{r - 2M} \right] \times \frac{[p - 2 - 2e \cos(\chi - w)][1 + e \cos(\chi - w)]^2}{M \sqrt{p^3 [(p - 2)^2 - 4e^2]}}$$



# Orbital evolution



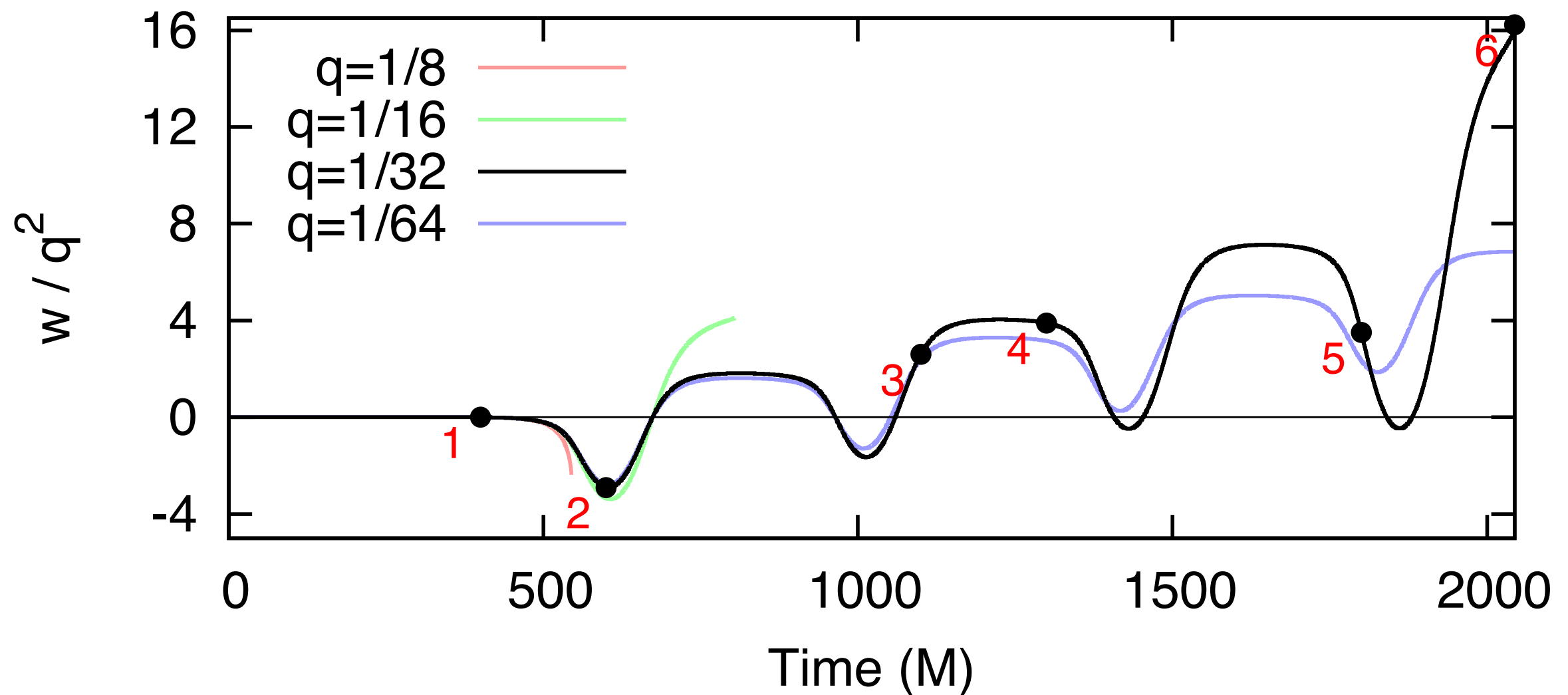
# Orbital evolution - “dissipative”



$$r(t) = \frac{Mp}{1 + e \cos(\chi - w)}$$



# Orbital evolution - “conservative”



$$r(t) = \frac{Mp}{1 + e \cos(\chi - w)}$$

# Waveforms at $\mathcal{I}^+$

