

The Chiral Anomaly

Some global symmetries in QFT are 'anomalous' in the quantum theory. eg Chiral symmetry in QCD

$$\mathcal{L} = i \bar{\Psi}_L \not{D} \Psi_L + i \bar{\Psi}_R \not{D} \Psi_R + m(\bar{\Psi}_L \Psi_R + \bar{\Psi}_R \Psi_L)$$

When $m=0$:
 N_f flavours

$$SU(N_f)_L \otimes SU(N_f)_R$$

$$\left(\begin{smallmatrix} 1 & & \\ & \ddots & \\ & & 1 \end{smallmatrix} \right) U(1)_L \otimes U(1)_R$$

$$\alpha = \beta : U(1)_B \quad \checkmark$$

$$\alpha = -\beta \quad U(1)_A \text{ is anomalous,}$$

$$\left. \begin{aligned} \Psi_L &\rightarrow U \Psi_L \\ \Psi_R &\rightarrow \tilde{U} \Psi_R \end{aligned} \right\} \text{these are fine...}$$

$$\Psi_L \rightarrow e^{i\alpha} \Psi_L$$

$$\Psi_R \rightarrow e^{i\beta} \Psi_R$$

$$U(1)_A: \quad \psi \rightarrow e^{i\alpha \gamma_5} \psi \quad \sim \begin{pmatrix} \psi_L \\ \psi_R \end{pmatrix} \rightarrow \begin{pmatrix} \psi_L \\ e^{i\alpha} \psi_R \end{pmatrix}$$

The breaking was first seen in one loop diagrams



$$\Sigma_{\sigma\nu\tau\rho} F^{\sigma\nu} F^{\tau\rho} \neq 0$$

If we tried to gauge $U(1)_A$ we would get a $A^\mu \gamma \gamma$ 3-pt vertex that is forbidden by gauge invariance.

Fujikawa's Non-Perturbative Proof

First I'm going to Fourier expand fields in path integral

$$\psi(x) = \sum_n a_n \phi_n(x)$$

$$\text{with } \phi_n(x) = u_r e^{-ikx}$$

$$\begin{matrix} \uparrow \\ \text{counts } r \\ \& k \end{matrix} \quad \begin{matrix} \uparrow \\ \end{matrix} \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} \dots \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

$$a_n = \int d^4x \bar{\phi}_n \psi$$

$$\bar{\psi}(x) = \sum_n \psi_n^\dagger(x) \bar{b}_n$$

$$\text{note orthogonality } \int \bar{\phi}_i \phi_j d^4x = \delta_{ij}$$

The key point is that $D[\psi] D[\bar{\psi}] \rightarrow \prod_n da_n \prod_m d\bar{b}_m$

$$\text{Under } U(1)_A \quad \psi \rightarrow \psi + i\omega \gamma_5 \psi + \dots$$

$$a_n \rightarrow a_n - i\omega \int d^4x \bar{\phi}_n \gamma_5 \psi \equiv a_n'$$

$\Pi da_n \rightarrow \Pi da_n'$ we get a det $\left| \frac{\partial \phi_i}{\partial a_n'} \right|$

$$\det \underset{\sim}{=} \det \begin{bmatrix} 1 & 0 & \dots & 0 \\ 0 & 1 & \dots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \dots & 1 \end{bmatrix}$$

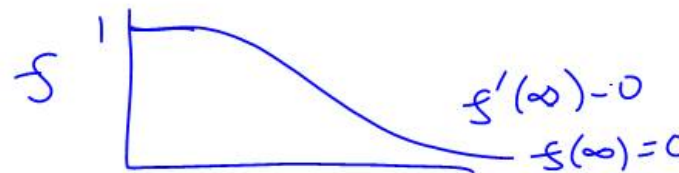
drop K_{ij} terms

since in det they square to small!

$$K_n = \omega \int \bar{\phi}_n \gamma_5 \phi_n d^4x$$

$$= 1 + i \sum_n K_n$$

$$\det \approx e^{+i \sum_n \omega \int \bar{\phi}_n \gamma_5 \phi_n d^4x}$$



Regulator: multiply by a smooth, Lorentz & gauge invariant g : $g\left[\frac{(i\not{D})^2}{g}\right]$

Now $\not{D}^2 = \gamma^\mu \gamma^\nu D_\mu D_\nu$

$$= \frac{1}{2} [\gamma^\mu, \gamma^\nu] D_\mu D_\nu + \frac{1}{2} \{\gamma^\mu, \gamma^\nu\} D_\mu D_\nu$$

$$= D^2 + \underbrace{\frac{1}{2} \gamma^\mu \gamma^\nu [D_\mu, D_\nu]}_{\frac{i}{2} \frac{1}{g} F_{\mu\nu} \gamma^\mu \gamma^\nu}$$

Taylor expand $\mathcal{F}(\frac{i \not{D}^2}{g})$

$$\begin{aligned} \mathcal{F}(\frac{i \not{D}^2}{g}) &= \mathcal{F}(-\frac{D^2}{g}) - \mathcal{F}'(-\frac{D^2}{g}) \frac{i}{2} F_{\mu\nu} \gamma^\mu \gamma^\nu \\ &\quad + \mathcal{F}''(-\frac{D^2}{g}) \left(-\frac{1}{8}\right) F_{\mu\nu} \gamma^\mu \gamma^\nu F_{\lambda\sigma} \gamma^\lambda \gamma^\sigma + \dots \end{aligned}$$

$$\det \simeq e^{i \sum_r \omega \int \bar{\psi}_n \gamma_5 f(-\not{\partial}_g^2) \psi_n d^4x}$$

$$\left\{ \sum_{r=1}^2 \int \frac{d^4k}{(2\pi)^4} \right\} \left\{ \bar{u}_r e^{-i k x} \right\} \left\{ u_r e^{i k x} \right\}$$

$$(1000) \begin{pmatrix} a_1 & a_2 & \dots \\ b_1 & b_2 & \dots \end{pmatrix} \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} = a_1$$

The sum over u_r gives $\text{Tr } M$

$$\simeq e^{i\omega \iint \frac{d^4k}{(2\pi)^4} d^4x \text{Tr } \gamma_5 f}$$

Now $\text{Tr } \gamma_5 = 0$, $\text{Tr } \gamma_5 \gamma^\mu \gamma^\nu = 0$, $\text{Tr } \gamma_5 \gamma^\mu \gamma^\nu \gamma^\lambda \gamma^\sigma = -4i \epsilon^{\mu\nu\lambda\sigma}$

$$\text{Tr } \gamma_5 \dots = 0 \text{ all higher}$$

$$\det = e^{i\omega \int \frac{d^4 k}{(2\pi)^4} f_x(-\frac{1}{8}) F_{\mu\nu} F_{\lambda\sigma} f'(-4i\epsilon^{\mu\nu\lambda\sigma})}$$

$$\text{Now } \int \frac{d^4 k}{(2\pi)^4} f''(-\frac{D^2}{9}) \xrightarrow[k]{\text{shift}} \int \frac{d^4 k'}{(2\pi)^4} f''(-k^2)$$

Wich rotate

$$= -\frac{i}{16\pi^2} \int k^2 dk^2 f''(k^2)$$

$$= -\frac{i}{16\pi^2} \left\{ [k^2 f']_0^\infty - \int_0^\infty f'(k^2) dk^2 \right\}$$

$$= -\frac{i}{16\pi^2} [f(0) - f(\infty)]$$

$$= -\frac{i}{16\pi^2}$$

Thus under U(1)_A

$$\Delta \mathcal{L} = +\frac{\omega}{32\pi^2} F \hat{F}$$

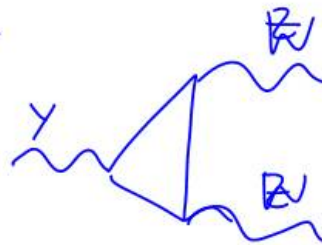
Also change $D[F]$

$$\rightarrow \Delta \mathcal{L} = +\frac{\omega}{16\pi^2} F \hat{F}$$

$\rightarrow$ this is not a symmetry!

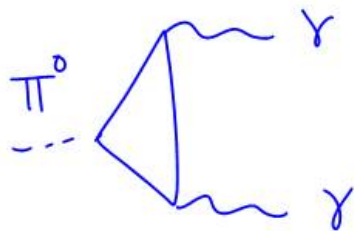
$U(1)_A \rightsquigarrow U(1)_X$ which is chiral.

You have to check



vanish.

They do because $\sum_f Y_f = 0$



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in Peskin Schroeder.