

NAGT

$$\mathcal{L}_{\text{NAGT}} = -\frac{1}{2} \text{Tr} F^\mu F_\mu + \bar{\Psi} (i \not{D} - m \mathbb{I}) \Psi$$

Gauge Fixing

Warm Up : consider the area integral $\int_0^L dx \int_0^L dy = L^2$

We can turn into a line integral by inserting $\delta(y)$

$$\int_0^L dx \int_0^L \delta(y) dy = \int_0^L dx = L$$

Move to polar (r, θ) coords

$\delta_y \delta_x \rightarrow \delta_\theta \delta_r$ in general get some //ogram in the new coords

$$\text{area} = \vec{a} \times \vec{b}$$


The area changes by a "Jacobian" factor (a determinant)

$$J = \begin{vmatrix} \frac{\partial x}{\partial r} & \frac{\partial x}{\partial \theta} \\ \frac{\partial y}{\partial r} & \frac{\partial y}{\partial \theta} \end{vmatrix} = \begin{vmatrix} \cos \theta & -r \sin \theta \\ r \sin \theta & r \cos \theta \end{vmatrix} = r$$

$$\iint dx dy \rightarrow \int J d\theta dr = \int r dr d\theta$$

Now to get a line integral along x insert $\delta(\sigma)$ (?)

$$\iint \delta(\sigma) r dr d\sigma = \int_0^L x dx = \frac{L}{2}$$

To correct this we need to change the normalization of δ -fn when change coords

$$\delta(y) \rightarrow \frac{1}{J|_{\sigma=\sigma_0}} \delta(\sigma - \sigma_0)$$

NAGT

Gauge freedom $A^\mu \rightarrow U A^\mu_T U^{-1} + \frac{i}{g} (d^\mu U) U^{-1}$

$\swarrow A^\mu T^a$ $\swarrow U T^a$ $\swarrow U^{-1} T^a$

For small ω^a $U = 1 + i \omega^a T^a + \dots$

$$\delta A^\mu_a = f^{abc} A^\mu_b \omega^c - \frac{1}{g} (d^\mu \omega^a)$$

This is a redundancy in the description ($\vec{E}, \vec{B}$ are gauge inv)

The overcounting causes trouble...

$$S = \int d^4x \overline{F^{a\mu\nu}} F^a_{\mu\nu} = \int d^4x \left(\overbrace{d^\mu A^\nu_a - d^\nu A^\mu_a}^{d^\mu A^\nu_a - d^\nu A^\mu_a} - \overbrace{d^\mu A^\nu_a d^\nu A^\mu_a}^{d^\mu A^\nu_a d^\nu A^\mu_a} - \overbrace{d^\nu A^\mu_a d^\mu A^\nu_a}^{d^\nu A^\mu_a d^\mu A^\nu_a} + \overbrace{d^\nu A^\mu_a d^\nu A^\mu_a}^{d^\nu A^\mu_a d^\nu A^\mu_a} \right)$$

$$\underline{\text{But}} \quad \int d^4x \quad \sum A_\mu^a \quad \underbrace{O_{ab}^{\mu\nu}}_{(g^{\mu\nu} \partial^2 - \partial^\mu \partial^\nu)} A_\nu^b$$

The propagator is the inverse of $\mathcal{O}$ but there isn't an inverse - I can tell because there is a zero eigenvalue

$$O^{\mu\nu} \partial_\nu \tilde{\omega}(\lambda) = 0$$

So we should restrict to solutions in some fixed gauge

$$D[A_\mu] \rightarrow D[A_\mu^{\omega=0}] D[\omega] \delta(\omega)$$

But we normally fix a gauge by a condition like

$$\partial^\mu A_\mu = \xi(x)$$

so we put in (?) $\delta[\partial^\mu A_\mu - \xi]$ for $\delta(w)$

we can so long as we include a "functional determinant" to take care of change of variables... $\mathcal{J} \leftarrow \det \frac{\delta(\partial^\mu A_\mu)}{\delta w} \Big|_{w=0}$

$$Z = \int D[A_\mu] \delta(\partial^\mu A_\mu - \xi) \mathcal{J} \Big|_{w=0} e^{iS}$$

Now $\delta(\partial^\mu A_\mu^a) = \xi^{abc} \partial^\mu A_\mu^b w^c - \frac{1}{g} \partial^\mu \partial_\mu w^a$

$$\begin{aligned}
 \Rightarrow \frac{\delta(\partial^\mu A_\mu^a)(x)}{\delta \omega^b(y)} &= f^{abc} \partial^\mu A_\mu^c \delta^4(x-y) - \frac{1}{g} \partial^\mu \partial_\mu \delta^4(x-y) \delta_b^a \\
 &= -\frac{1}{g} \delta^4(x-y) \partial^\mu D_\mu^{ab}(x) \\
 &\quad \uparrow \partial_\mu \delta^{ab} - g f^{abc} A_{c\mu}
 \end{aligned}$$

$\uparrow$
 generators
 of adjoint
 rep.

We want to take the determinant

$$\det M = \int D[\eta] D[\bar{\eta}] e^{-\int d^4x \bar{\eta}(x) M \eta(x)}$$

$\uparrow$ Grassmannian

η & $\bar{\eta}$ are "Faddeev-Popov" ghosts - not physical but occur in loops in Feynman diagrams

They are scalars (they get a - sign on loops)

$$\frac{1}{p^2}$$

$$g f^{abc} q^m$$

$$\int d^4x \mathcal{L} e^{iS}$$

One final trick.. any choice of $f(x)$ gives the same answer.. pick two, add & you'll get twice Z ... why not pick a smeared distribution of f s?

$$Z = \int D[f] e^{i \int d^4x \frac{1}{2(1-\xi)} f^2(x)} \int D[A_\mu] \delta(\partial A - f) \int e^{iS}$$

do ξ integral.

$$Z = \int D[A_\mu] J e^{iS} e^{i \int d^4x \frac{1}{2(1-\xi)} (\partial^\mu A_\mu)^2}$$

↳ integrate by parts
 $- A_\mu \partial^\mu \partial^\nu A_\nu$

Now $O_{ab}^{\mu\nu} = i \delta_{ab} \left(g^{\mu\nu} \partial^2 + \frac{\xi}{1-\xi} \partial^\mu \partial^\nu \right)$

The inverse is $- \delta_{ab} \left(g_{\mu\nu} - \frac{\xi p_\mu p_\nu}{p^2} \right) \frac{i}{p^2}$